

MAXIMAL SUPERSYMMETRY

"30 years of supergravity"  
ENS, Paris, 16-20 October

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(AEI, Golm)

Symmetry = (arguably) most successful  
guiding principle of physics

⇒ search for maximally symmetric theories still best strategy for unification of fundamental interactions & consistent quantization of gravity

Supersymmetry (SUSY) = enlargement of conventional (= Lie algebra) notion of symmetry  
→ B. Zumino  
J. Wess

Local SUSY ↔ Supergravity (SUGRA)

→ P. van Nieuwenhuizen

SUSY algebras exist for all

$N = \#$  (supercharges)

$D = \#$  (dimensions of space-time)

However, consistent and interacting SUSY field theories appear to exist only for limited values

$$N \leq 32 \iff D \leq 11$$

[for global SUSY :  $N \leq 16 \iff D \leq 10$ ]

Search for maximally supersymmetric theories has led to maximal point FT's

\*  $D = 11$  SUGRA (CJS, 1978)

\*  $D = 4, N = 8$  SUGRA

(CJ, 1979; DMN, 1982)

as well as maximally supersymmetric extended object theories:

\* Type II superstrings (GS, 1982)

→ M. Green, J. Schwarz

\*  $D = 11$  supermembrane (BST, 1987)

↕ (DMHN, 1988)

M(atrix) Theory  $\equiv$   $U \rightarrow \infty$  SUSY YM

which contain  $D = 10$  &  $D = 11$  SUGRAS as low energy limits (in some sense)

Remarkable fact: there are distinguished theories in the "space of all theories"!

Drawback: rigid structure of these

theories makes it difficult to establish link with the "real world" ...

Bounds on  $N$  and  $D$  may be related to:

- absence of massless spin  $> 2$

(Nahm, 1976)

- connection with maximally extended Kac-Moody-Lie algebras  $E_8$  and  $E_{10}$

→ B. Julia, P. West

Important Caveat: all constructions

based on "conventional" notions of space-time, but quantum gravity very likely requires more than a space-time based QFT →

- non-commutative space-time, quantum geometry & symmetry?
- emergent space-time?

[e.g. as suggested by link between BKL cosmology and  $E_{10}$  (DHN)]

- other ideas?

→ may be relevant also for

question of SUSY breaking!



# D = 11 : the mother of all SUGRAS?

(CJS, 1978)

Field content:  $G_{MN}$  (44),  $A_{MNP}$  (84),  $\psi_M$  (128)

Lagrangian:

$$E^{-1} \mathcal{L} = \frac{1}{4} R - \frac{i}{2} \bar{\psi}_M \Gamma_{MNP} \mathcal{D}_N \psi_P - \frac{1}{4} F_{MNPQ}^2$$

$$- \frac{i}{96} \left( \bar{\psi}_M \Gamma_{MNPQRS} \psi_N + 12 \bar{\psi}_P \Gamma_{PQRS} \psi^S \right) F_{PQRS}$$

$$+ \frac{1}{2} \epsilon_{M_1 \dots M_{11}} F_{M_1 M_2 M_3 M_4} F_{M_5 M_6 M_7 M_8} A_{M_9 M_{10} M_{11}}$$

+ quartic terms

Unique theory:

\* no matter couplings

\* not deformable

\* maximal space-time dimension

compatible with SUSY

\* multiplet structure "explained" by  $E_{11}$

Furthermore:

\* compactifies preferentially via KK

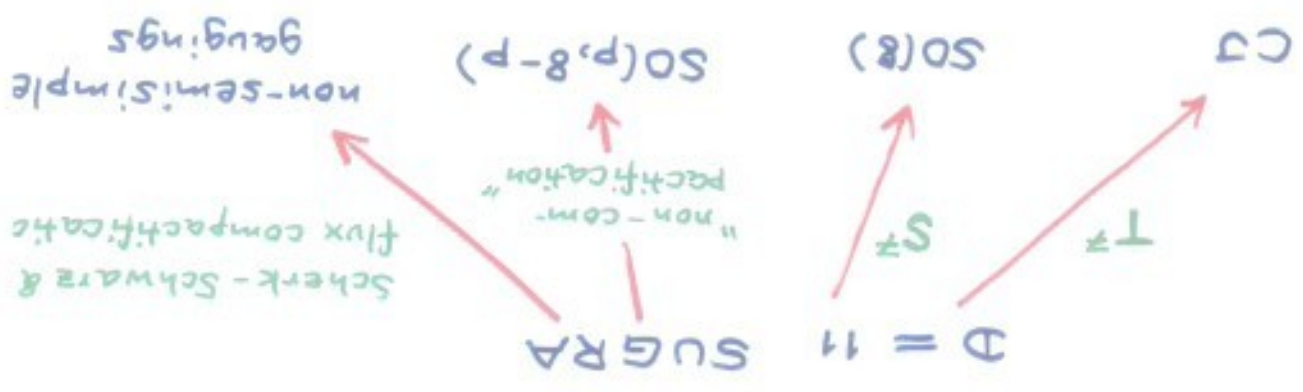
to  $D = 4$  (and  $D = 7$ ) (FR, 1980)

← M. Duff

\* strong coupling limit of

superstring theory? (Witten, 1995)

Kaluza Klein compactification to  $D=4$ :



→ S. Ferrara

Also contains ungauged maximal SUGRAS  
 $D$  via dimensional reduction on  $T^{11-D}$ .

Idem:  $D=11$  SUGRA on  $AdS^7 \times S^4$

$D=11$  SUGRA on  $AdS^5 \times S^5$

However, not all maximal theories  
in  $D < 11$  have higher dimensional origin

→ Richest variety of maximal gauged

SUGRAS for  $D=3$ : (Samtleben, N., 2000)

\* deformations of  $D=3, N=16$

without  $D > 3$  origin

\* numerous semisimple and

non-semisimple  $G_0 \subset E_{8(8)}$

\* KK compactifications to  $D=3$

↔ non-semisimple  $G_0$ , e.g.

$IIA$  on  $T^7 \leftrightarrow G_0 = SO(8) \times T^{28}$

## Maximal $\mathcal{D}=3$ SUGRAS with semisimple $G_0$

gauge group $G_0$	ratio $g_1/g_2$	$(n_L, n_R)$	ground state symmetry group
$SO(8) \times SO(8)$	$g_1/g_2 = -1$	$(8, 8)$	$OSP(8 2, \mathbb{R}) \times OSP(8 2, \mathbb{R})$
$SO(7, 1) \times SO(7, 1)$	$g_1/g_2 = -1$	$(8, 8)$	$F(4) \times F(4)$
$SO(6, 2) \times SO(6, 2)$	$g_1/g_2 = -1$	$(8, 8)$	$SU(4 1, 1) \times SU(4 1, 1)$
$SO(5, 3) \times SO(5, 3)$	$g_1/g_2 = -1$	$(8, 8)$	$OSP(4^* 4) \times OSP(4^* 4)$
$SO(4, 4) \times SO(4, 4)$	$g_1/g_2 = -1$	$(8, 8)$	Minkowski vacuum
$G_{2(2)} \times F_{4(4)}$	$g_{02}/g_{r4} = -3/2$	$(4, 12)$	$D^1(2, 1; -\frac{3}{2}) \times OSP(4^* 6)$
$G_2 \times F_{4(-20)}$	$g_{02}/g_{r4} = -3/2$	$(7, 9)$	$G(3) \times OSP(9 2, \mathbb{R})$
$E_{6(6)} \times SL(3)$	$g_{v2}/g_{r6} = -2$	$(16, 0)$	$OSP(4^* 8) \times SU(1, 1)$
$E_{6(2)} \times SU(2, 1)$	$g_{v2}/g_{r6} = -2$	$(12, 4)$	$SU(6 1, 1) \times D^1(2, 1; -\frac{2}{3})$
$E_{6(-14)} \times SU(3)$	$g_{v2}/g_{r6} = -2$	$(10, 6)$	$OSP(10 2, \mathbb{R}) \times SU(3 1, 1)$
$E_{7(7)} \times SL(2)$	$g_{v1}/g_{r7} = -3$	$(16, 0)$	$SU(8 1, 1) \times SU(1, 1)$
$E_{7(-5)} \times SU(2)$	$g_{v1}/g_{r7} = -3$	$(12, 4)$	$OSP(12 2, \mathbb{R}) \times D^1(2, 1; -\frac{3}{2})$
$E_{8(8)}$	$g_{r8}$	$(16, 0)$	$OSP(16 2, \mathbb{R}) \times SU(1, 1)$

Table 2: The  $N = 16$  theories with semisimple gauge groups  $G_0$ . Except for the last row, the gauge groups appear as direct products of two factors whose coupling constant ratio  $g_1/g_2$  is determined by (9.2). All these theories admit a maximally supersymmetric AdS (or Minkowski, for  $G_0 = SO(4, 4) \times SO(4, 4)$ ) ground state, whose symmetry group factorizes according to  $G_L \times G_R$ , as specified in the last column; the supercharges split accordingly into  $n_L + n_R = 16$ .

as well as  $G_0 = SO(8, \mathbb{C}) \subset E_{8(8)}$ :  
 no isometries, unstable de Sitter!

Further peculiarities of  $\mathcal{D} = 3$ :

\* Fully SUSY (hence, stable) vacua

also for non-compact gaugings

\*  $\exists$  stable AdS ground states  
 with fully broken SUSY



Summary: fairly complete picture of bestiary of maximal SUSY theories, but several "old" problems persist:

\* Off-shell formulation (components or superspace) of maximal theories?

\* Finiteness (or not) of  $N = 8$  SUGRA? [NB:  $D = 11$  SUGRA is divergent!]

→ P. Townsend, S. Deser

\* Maximal SUSY and higher order corrections ( $R^2, R^3, \dots$ )?

\*  $E_{n(n)}$  and higher order corrections? (e.g.  $E_{n(n)} \rightarrow E_n(\mathbb{Z})$ ?)

and, perhaps, less importantly:

\* Extremal structure of gauged SUGRA potentials:

•  $N = 8, D = 4$ : no progress since Warner ('88)

•  $N = 16, D = 3$ : much more complicated

e.g.  $G_0 = G_2 \times F_{4(-20)}$  on subspace of  $SU(3) \times SU(3)$  singlets

(T. Fischbacher, PhD Thesis)







$$\begin{aligned}
&= \frac{24125}{2048} + \frac{9}{2048} \cos(8\tau_5) + \frac{9}{2048} \cos(8\tau_2) \\
&+ \frac{119}{128} \sin(4\tau_5) \sin(4\tau_2) - \frac{27}{2048} \cos(8\tau_2) \cos(8\tau_5) \\
&- \frac{1}{64} \cos(4\tau_2) \cos(4\tau_3 - 4\tau_6) \cos(4\tau_5) - \frac{9}{2048} \cos(8\tau_1 - 8\tau_4) \\
&- \frac{9}{2048} \cos(8\tau_1 - 8\tau_4) \cos(8\tau_5) - \frac{1}{64} \cos(4\tau_1 - 4\tau_4) \cos(4\tau_3 - 4\tau_6) \\
&+ \frac{1}{64} \sin(4\tau_5) \sin(4\tau_3 - 4\tau_6) \sin(4\tau_1 - 4\tau_4) - \frac{9}{2048} \cos(8\tau_1 - 8\tau_4) \cos(8\tau_2) \\
&- \frac{517}{128} \cos(4\tau_1 - 4\tau_4) \sin(8\tau_5) \sin(8\tau_2) \\
&+ \frac{115}{128} \cos(4\tau_1 - 4\tau_4) \cos(4\tau_2) \cos(4\tau_5) \\
&- \frac{9}{2048} \cos(8\tau_1 - 8\tau_4) \cos(8\tau_2) \cos(8\tau_5) \\
&+ \frac{1}{64} \sin(4\tau_3 - 4\tau_6) \sin(4\tau_2) \sin(4\tau_1 - 4\tau_4) \\
&- \frac{1}{64} \cos(4\tau_1 - 4\tau_4) \cos(4\tau_3 - 4\tau_6) \sin(4\tau_5) \sin(4\tau_2) + \frac{449}{512} \cosh(z) \\
&- \frac{65}{2048} \cosh(2z) - \frac{3}{512} \cosh(z) \cos(8\tau_5) + \frac{3}{2048} \cosh(2z) \cos(8\tau_5) \\
&- \frac{512}{512} \cosh(z) \cos(8\tau_2) + \frac{3}{2048} \cosh(2z) \cos(8\tau_2) \\
&- \frac{29}{32} \cosh(z) \sin(4\tau_5) \sin(4\tau_2) + \frac{126}{128} \cosh(2z) \sin(4\tau_5) \sin(4\tau_2) \\
&+ \frac{512}{512} \cosh(z) \cos(8\tau_2) \cos(8\tau_5) - \frac{9}{2048} \cosh(2z) \cos(8\tau_2) \cos(8\tau_5) \\
&+ \frac{1}{64} \cosh(2z) \cos(4\tau_2) \cos(4\tau_3 - 4\tau_6) \cos(4\tau_5) + \frac{512}{512} \cosh(z) \cos(8\tau_1 - 8\tau_4) \\
&- \frac{3}{2048} \cosh(2z) \cos(8\tau_1 - 8\tau_4) + \frac{3}{512} \cosh(z) \cos(8\tau_1 - 8\tau_4) \cos(8\tau_5) \\
&- \frac{3}{2048} \cosh(2z) \cos(8\tau_1 - 8\tau_4) \cos(8\tau_5) \\
&+ \frac{1}{64} \cosh(2z) \cos(4\tau_1 - 4\tau_4) \cos(4\tau_3 - 4\tau_6) \\
&- \frac{1}{64} \cosh(2z) \sin(4\tau_5) \sin(4\tau_3 - 4\tau_6) \sin(4\tau_1 - 4\tau_4) \\
&+ \frac{3}{512} \cosh(z) \cos(8\tau_1 - 8\tau_4) \cos(8\tau_2) \\
&- \frac{3}{2048} \cosh(2z) \cos(8\tau_1 - 8\tau_4) \cos(8\tau_2) \\
&+ \frac{1}{128} \cosh(z) \cos(4\tau_1 - 4\tau_4) \sin(8\tau_5) \sin(8\tau_2) \\
&- \frac{3}{512} \cosh(2z) \cos(4\tau_1 - 4\tau_4) \sin(8\tau_5) \sin(8\tau_2) \\
&- \frac{29}{32} \cosh(z) \cos(4\tau_1 - 4\tau_4) \cos(4\tau_2) \cos(4\tau_5) \\
&+ \frac{512}{512} \cosh(z) \cos(8\tau_1 - 8\tau_4) \cos(8\tau_2) \cos(8\tau_5) \\
&+ \frac{3}{512} \cosh(z) \cos(8\tau_1 - 8\tau_4) \cos(4\tau_2) \cos(4\tau_5) \\
&+ \frac{126}{128} \cosh(2z) \cos(8\tau_1 - 8\tau_4) \cos(8\tau_2) \cos(4\tau_5) \\
&- \frac{3}{2048} \cosh(2z) \cos(8\tau_1 - 8\tau_4) \cos(8\tau_2) \cos(8\tau_5) \\
&- \frac{1}{64} \cosh(2z) \sin(4\tau_3 - 4\tau_6) \sin(4\tau_2) \sin(4\tau_1 - 4\tau_4) \\
&+ \frac{1}{64} \cosh(2z) \cos(4\tau_1 - 4\tau_4) \cos(4\tau_3 - 4\tau_6) \sin(4\tau_5) \sin(4\tau_2) \\
&+ \frac{449}{512} \cosh(s) - \frac{65}{2048} \cosh(2s) - \frac{3}{512} \cosh(s) \cos(8\tau_5) \\
&+ \frac{3}{2048} \cosh(2s) \cos(8\tau_5) - \frac{517}{128} \cosh(s) \cos(8\tau_2) \\
&+ \frac{3}{2048} \cosh(2s) \cos(8\tau_2) - \frac{3}{512} \cosh(s) \cos(8\tau_2) \\
&+ \frac{126}{128} \cosh(2s) \sin(4\tau_5) \sin(4\tau_2) + \frac{9}{512} \cosh(s) \cos(8\tau_2) \cos(8\tau_5) \\
&- \frac{9}{2048} \cosh(2s) \cos(8\tau_2) \cos(8\tau_5) \\
&+ \frac{1}{64} \cosh(2s) \cos(4\tau_2) \cos(4\tau_3 - 4\tau_6) \cos(4\tau_5) \\
&+ \frac{3}{512} \cosh(s) \cos(8\tau_1 - 8\tau_4) - \frac{3}{2048} \cosh(2s) \cos(8\tau_1 - 8\tau_4) \\
&\dots
\end{aligned}$$

$$\begin{aligned}
&+ \frac{1}{256} \sin(2\tau_3 - 2\tau_6) \sin(2\tau_2 + 2\tau_5) \sin(6\tau_1 - 6\tau_4) \sinh(2z) \sinh(2s) \\
&- \frac{1}{256} \sin(2\tau_3 - 2\tau_6) \sin(6\tau_2 - 6\tau_5) \sin(6\tau_1 - 6\tau_4) \sinh(2z) \sinh(2s) \\
&- \frac{1}{256} \sin(2\tau_3 - 2\tau_6) \sin(6\tau_2 + 6\tau_5) \sin(6\tau_1 - 6\tau_4) \sinh(2z) \sinh(2s) \\
&- \frac{1229}{64} \cos(2\tau_1 - 2\tau_4) \cos(2\tau_2 - 2\tau_5) \cos(2\tau_3 - 2\tau_6) \sinh(z) \sinh(s) \\
&- \frac{1}{64} \cos(2\tau_1 - 2\tau_4) \cos(2\tau_2 + 6\tau_5) \cos(2\tau_3 - 2\tau_6) \sinh(z) \sinh(s) \\
&+ \frac{3}{64} \cos(2\tau_1 - 2\tau_4) \cos(6\tau_2 - 6\tau_5) \cos(2\tau_3 - 2\tau_6) \sinh(z) \sinh(s) \\
&- \frac{1}{64} \cos(2\tau_1 - 2\tau_4) \cos(6\tau_2 + 6\tau_5) \cos(2\tau_3 - 2\tau_6) \sinh(z) \sinh(s) \\
&+ \frac{1}{64} \cos(6\tau_1 - 6\tau_4) \cos(6\tau_2 + 2\tau_5) \cos(2\tau_3 - 2\tau_6) \sinh(z) \sinh(s) \\
&+ \frac{1}{64} \cos(6\tau_1 - 6\tau_4) \cos(6\tau_2 - 6\tau_5) \cos(2\tau_3 - 2\tau_6) \sinh(z) \sinh(s) \\
&+ \frac{1}{64} \cos(6\tau_1 - 6\tau_4) \cos(6\tau_2 + 2\tau_5) \cos(2\tau_3 - 2\tau_6) \sinh(z) \sinh(s) \\
&+ \frac{1}{64} \cos(6\tau_1 - 6\tau_4) \cos(6\tau_2 - 6\tau_5) \cos(2\tau_3 - 2\tau_6) \sinh(z) \sinh(s) \\
&+ \frac{1}{64} \cos(6\tau_1 - 6\tau_4) \cos(6\tau_2 + 6\tau_5) \cos(2\tau_3 - 2\tau_6) \sinh(z) \sinh(s) \\
&+ \frac{1}{64} \cos(6\tau_1 - 6\tau_4) \cos(6\tau_2 - 6\tau_5) \cos(2\tau_3 - 2\tau_6) \sinh(z) \sinh(s) \\
&+ \frac{1}{64} \cos(6\tau_1 - 6\tau_4) \cos(6\tau_2 + 2\tau_5) \cos(2\tau_3 - 2\tau_6) \sinh(z) \sinh(s) \\
&+ \frac{1}{64} \cos(6\tau_1 - 6\tau_4) \cos(6\tau_2 - 6\tau_5) \cos(2\tau_3 - 2\tau_6) \sinh(z) \sinh(s) \\
&- \frac{1}{64} \cosh(2s) \cosh(2z) \cos(4\tau_1 - 4\tau_4) \cos(4\tau_3 - 4\tau_6) \sinh(4\tau_5) \sinh(4\tau_2)
\end{aligned}$$

## General pattern:

R symmetry = maximal compact subgroup  $K(E_n) \subset E_n$

Bosons  $\in$  single valued reps

Fermions  $\in$  double valued (spinorial) reps

$$D=4 : K(E_7) = SU(8) \subset E_{7(7)}$$

$$\text{Bosons : } 1 \oplus 28 \oplus 70$$

$$\text{Fermions : } 8 \oplus 56$$

$$\rightarrow SU(8)/\mathbb{Z}_2$$

$$D=3 : K(E_8) = "SO(16)" \subset E_{8(8)}$$

$$\text{Bosons : } 128_s$$

$$\text{Fermions : } 16_v \oplus 128_c$$

$$\text{NB : center of Spin}(16) = \mathbb{Z}_2^L \times \mathbb{Z}_2^R \rightarrow "SO(16)" = Spin(16)/\mathbb{Z}_2^L$$

(Klein 2-folds)

Does this extend to  $K(E_9)$  and  $K(E_{10})$ ?



Fermions and maximal SUSY in  $D=2$

$K(E_g) :=$  subgroup of  $E_{g(g)}$  left invariant by generalized **Cherny** involution

in terms of loop group realization:

$$u(\lambda) \in K(E_g) \iff u(\lambda) = [(u^T)^{-1}] \left( \frac{1}{\lambda} \right)$$

$\downarrow$ 
 $\downarrow$

$\in E_g(\mathbb{C})$ 
 $\in E_g(\mathbb{C})$

for all  $\lambda \in \mathbb{C}$

$\Rightarrow$  at fixed points  $\lambda = \frac{1}{\lambda} = \pm 1 \rightarrow SO(16)_{\pm}$

\*  $K(E_g)$  is not Kac-Moody (and  $\neq \widehat{SO(16)}$ )

\*  $K(E_g)$  is not simple (unlike  $E_{g(g)}$ ):

$\exists$  non-trivial ideals related to unfaithful spinorial representations

More specifically,  $K(E_g)$  acts on chiral

components of fermions via the "evaluation map" (Samtleben, N.)

$$\chi_{\pm}'(x) = u(x, \lambda) \Big|_{\lambda = \pm 1} \chi_{\pm}(x)$$

$\downarrow$  128c

looks like "generalized holonomy" group  $SO(16)_{+} \times SO(16)_{-} \dots$

... but action on full set of SUGRA fermions (and their derivatives) is more complicated, as it involves

$$U^{(n)}(r) \Big|_{r \neq 1} \quad \text{for } n \geq 1$$

Few results and many questions:

- \*  $K(E_g)$  as a generalized holonomy
- \* Finite dimensional "generalized holonomy groups" from quotient algebras  $K(E_g)/\mathfrak{g}$  such that e.g.  $\mathfrak{g} = \mathfrak{g}_{\text{Dirac}} \leftrightarrow SO(16)_+ \times SO(16)_-$
- \* Faithful (hence,  $\infty$ -dimensional) spinorial representations of  $K(E_g)$ ?
- \*  $K(E_g)$  as a generalized **R symmetry**: does there exist an  $\infty$ -dimensional SUSY algebra with  $K(E_g)$ ?

idem for  $K(E_{10}) \subset E_{10}$

... but much harder!

→ DKN, DBHP

Is there path leading from top to bottom, or vice versa?

the real (non-SUSY) world



Maximal SUSY theories



Identify

$$SU(3) = [SU(3)_c \times SU(3)_f]_{diag}$$

(M. Gell-Mann, '83)

$$\begin{pmatrix} u \\ c \\ t \end{pmatrix}_L$$

$$\begin{pmatrix} t \\ n \\ s \end{pmatrix}_L$$

$$\begin{pmatrix} d \\ s \\ b \end{pmatrix}_L$$

$$\begin{pmatrix} \bar{u} \\ \bar{d} \\ \bar{s} \end{pmatrix}_L$$

$$\begin{pmatrix} \bar{u} \\ \bar{c} \\ \bar{t} \end{pmatrix}_L$$

$$\begin{pmatrix} \bar{u} \\ \bar{d} \\ \bar{s} \end{pmatrix}_L$$

$$\begin{pmatrix} \bar{u} \\ \bar{c} \\ \bar{t} \end{pmatrix}_L$$

$$\begin{pmatrix} \bar{u} \\ \bar{d} \\ \bar{s} \end{pmatrix}_L$$

$$3_c \times 3_f \rightarrow 8 + 1$$

$$3_c \times 3_f \rightarrow 8 + 1$$

$$3_c \times 3_f \rightarrow 6 + \bar{3}$$

$$3_c \times 3_f \rightarrow 6 + \bar{3}$$

$$3_c \times 3_f \rightarrow 3$$

$$3_c \times 3_f \rightarrow 3$$

$$3_c \times 3_f \rightarrow 3$$

$$3_c \times 3_f \rightarrow 3$$

$$Q = \frac{2}{3} - 9$$

$$Q = -\frac{2}{3} + 9$$

$$Q = -\frac{1}{3} + 9$$

$$Q = \frac{1}{3} - 9$$

$$Q = -1 + 9$$

$$Q = 1 - 9$$

$$Q = -9$$

$$Q = +9$$

representations match with  $N = 8$  for spinon charge  $q = \frac{6}{1}!$

\* Scheme is realized at  $SU(3) \times U(1)$

Stationary point of  $N = 8$  SUGRA

(Warner, N., 1985)

... but:

... a mirage?



\* Number of fermionic states in

$$\text{nature is } 3 \times 16 = 48 = 56 - 8 \text{ (so far...)}$$

\* Unbroken symmetry of low

energy physics is  $SU(3) \times U(1)$

\* Analogous to Wilczek's scheme

of color-flavor-locking

→ hep-ph/0003183

Remember Albert Einstein:

"The Lord is subtle, but not malicious"