

TOPICS IN WALL-CROSSING

Note Title

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"WALL-CROSSING FROM SUSY GALAXIES" 1008.0030

"BOUND-STATE TMN WALLS" 1008.3555

GAIUTTO - M - NEITZKE

"4D WC + 3D FT" 0807.4723

"WC, HITCHIN, WKB" 0907.3987

"FRAMED BPS STATES" 1006.0146

WALL-CROSSING
IN COUPLED 2D-4D ????? . ?????

Review: PITP LECTURES
ON BPS STATES....

Homepage

I. INTRODUCTION

WALL-CROSSING IN $D=4, N=2$ THEORIES HAS BEEN KNOWN SINCE THE WORK OF SEIBERG & WITTEN.

THE 4D PHENOMENA WAS INSPIRED BY THE DISCOVERY OF A SIMILAR PHENOMENA IN THEORIES WITH $D=2$ $(2,2)$ SUSY (F, C, I, V) .

CECOTTI & VAFA WENT ON TO GIVE W.C.F. FOR THESE $D=2$ SYSTEMS, BUT ONLY IN THE PAST FEW YEARS HAVE WE FOUND QUANTITATIVE W.C.F. FOR $D=4, N=2$ SYSTEMS.

THIS TALK: AN OVERVIEW OF A POINT OF VIEW I HAVE DEVELOPED WITH MY COLLABORATORS:

MAIN RESULTS OF THIS OVERVIEW:

1. NEW INDICES ξ & NEW BPS STATES:

PSC ξ & FRAMED BPS STATES

2. HALO PICTURE OF WALL CROSSING AND A NEW DERIVATION OF THE KSWCF VIA "BPS GALAXIES"

3. LINE DEFECTS ξ & FRAMED BPS STATES LEAD TO "MOTIVIC KSWCF" AND Q-DEFORMATION OF FUNCTIONS ON SEIBERG-WITTEN MODULI SPACES.

4. LINE DEFECTS \Rightarrow DARBOUX COORDINATES \Rightarrow CONSTRUCTION OF HK METRICS ON \mathcal{M}

5. FOR "A₁-THEORIES" DARBOUX COORD'S COINCIDE WITH FOCK-GONCHAROV COORD'S

6. SURFACE DEFECTS \Rightarrow GENERALIZED 2D-4D WCF.

OVERVIEW OF THE OVERVIEW

II. $\mathcal{N}=2$ REVIEW

- (A.) BPS STATES & THEIR INDICES
- (B.) SELF-DUAL ABELIAN GAUGE THEORY
- (C.) MARGINAL STABILITY & THE WALL-CROSSING PROBLEM.

III. PRIMITIVE & SEMI-PRIMITIVE WCF

- (A.) MULTI-CENTER SOLUTIONS
- (B.) BOUNDSTATE RADIUS
- (C.) BASIC WC-MECHANISM
- (D.) AMS PARADOX
- (E.) HALOS + HALO FOCK SPACE
- (F.) SEMIPRIMITIVE WCF

IV. BPS GALAXIES $\frac{1}{\epsilon}$ KSWCF

(A.) NEED TO GO BEYOND SEMI-PRIMITIVE

(B.) BPS GALAXIES
- ENTROPIC SUPPRESSION

(C.) FRAMED BPS STATES, $\bar{\Omega}$

(D.) FRAMED BPS WC.: THE
KS-TRANSFORMATIONS

(E.) DERIVING THE KSWCF

(F.) RESOLVING AMS PARADOX:
CONJUGATION WALLS $\frac{1}{\epsilon}$,
GENERALIZATION OF KSWCF

V. LINE DEFECTS $\frac{1}{\ell}$ FRAMED BPS STATES

(A.) DEF. OF $L_{\mathcal{L}}$

$$(B.) \mathcal{H}_{L_{\mathcal{L}}} = \bigoplus_{\gamma} \mathcal{H}_{L_{\mathcal{L}}, \gamma}$$

FRAMED BPS STATES

$$E \geq -\operatorname{Re}(z_{\gamma}/\ell)$$

(C.) GENERATING FUNCTION
HALOS AGAIN

MOTIVIC KSWCF

(D.) STRONG POSITIVITY,
CLUSTER ALGEBRAS,
AND TROPICAL LABELS

VI : HYPERKÄHLER GEOMETRY:

COMPACTIFICATION TO THREE DIMENSIONS

- (A.) \mathcal{M} AS TORUS FIBRATION
- (B.) DARBOUX COORDINATE EXPANSION OF $\langle L_5 \rangle$
- (C.) PROPERTIES OF \mathcal{Y}_γ
- (D.) CONSTRUCTING THE \mathcal{Y}_γ VIA TBA
- (E.) CONSTRUCTION OF HK METRICS
- (F.) QUANTUM DEFORMATION OF \mathcal{M}

VII. M5 BRANES & HITCHIN SYSTEMS

(A.) $\Delta = (2,0)$ $\mathcal{N} = 2$ THEORIES / $C \times \mathbb{R}^{1,3}$

(B.) 3D COMPACT: $\mathcal{M} =$ HITCHIN MODULI

(C.) \exists \hookrightarrow SET OF LINE OPERATORS
LABELED BY $\rho \in C$

(D.) $\langle L_S(\rho) \rangle = \text{TrHol}$ (FLAT GAUGE FLD)

(E.) FOR A_1 -HITCHIN SYSTEMS

- $\mathcal{Y}_\gamma =$ FOCK-GONCHAROV COORD'S
- \exists ALGORITHM FOR COMPUTING
 - BPS INDICES $\Omega(\gamma)$
 - FRAMED INDICES $\overline{\Omega}(\gamma)$

VIII. SURFACE DEFECTS $\frac{1}{2}$ 2D/4D WCF

(A.) DEFINITION OF SURFACE DEFECTS
PRESERVED SUSY

(B.) EXAMPLES

(C.) IR DESCRIPTION :
GUKOV-WITTEN PARAMETERS
 $\frac{1}{2}$ TORSORS \mathbb{T}_i

(D.) INCLUDING LINE DEFECTS

(E.) FORMAL STRUCTURE OF
2D/4D WCF

(F.) 3D GEOMETRY: HHVB'S
 $V \rightarrow M$

(G.) EXAMPLES: HITCHIN SYSTEMS

II. $\mathcal{N}=2$ REVIEW

STUDY PARTICLE SPECTRUM
OF THEORIES WITH $D=4, \mathcal{N}=2$:

(A.) SUSY ALGEBRA

$$\mathcal{A} = \mathcal{A}^0 \oplus \mathcal{A}^1$$

$$\mathcal{A}^0 = \text{spin}(1,3) \oplus \text{su}(2)_R \oplus \text{u}(1)_R \oplus \mathbb{C}$$

$$\mathcal{A}^1 = \left[(2, 1; 2)_{+1} \oplus (1, 2; 2)_{-1} \right]_R$$

$$\left(Q_\alpha^A \right)^\dagger = \bar{Q}_{\dot{\alpha}A}$$

$$\{ Q_\alpha^A, \bar{Q}_{\dot{\beta}B} \} = 2 \sigma_{\alpha\dot{\beta}}^m P_m \delta^A_B$$

$$\{ Q_\alpha^A, Q_\beta^B \} = 2 \epsilon_{\alpha\beta} \epsilon^{AB} \bar{Z}$$

(B.) A USEFUL INVOLUTION

FOR ANY PHASE \mathcal{S} DEFINE
AN INVOLUTION OF $SU(2,2|2)$

BY $\vec{X} \rightarrow -\vec{X} \quad \& \quad U(1)_R$ ROTATION BY \mathcal{S}

$$\mathcal{J} = \mathcal{J}^{'+} \oplus \mathcal{J}^{'-}$$

$$R_{\alpha}^A := \mathcal{S}^{-1} Q_{\alpha}^A + \mathcal{S} \sigma_{\alpha\beta}^0 \bar{Q}^{\beta A}$$

$$J_{\alpha}^A := \mathcal{S}^{-1} Q_{\alpha}^A - \mathcal{S} \sigma_{\alpha\beta}^0 \bar{Q}^{\beta A}$$

$$\mathcal{S} = \sqrt{\mathcal{S}}$$

(C.) REPRESENTATIONS

MASSIVE PARTICLE = UNITARY IRREP.
INDUCED FROM F.D.
REP OF $\Delta_l^0 \oplus \Delta_l^1$

$$\Delta_l^0 = \mathfrak{so}(3) \oplus \mathfrak{su}(2)_R \oplus \mathfrak{u}(1)_R \oplus \mathbb{C}$$

$$\left(\mathcal{R}_l + (\mathcal{R}_l)^+ \right)^2 = 4(E + \operatorname{Re}(Z/\zeta))$$

$$\Rightarrow \boxed{E \geq -\operatorname{Re}(Z/\zeta)}$$

c.f. LINE DEFECTS BELOW

FOR $Z = e^{i\alpha} |Z|$ CHOOSE

$$\zeta = -e^{i\alpha} \Rightarrow \boxed{M \geq |Z|}$$

$\mathbb{R}_\alpha^A \cong \mathbb{J}_\alpha^A$ FORM

(GRADED COMMUTING) CLIFFORD ALGEBRAS. EACH HAS IRREP

$$\rho_{hh} = (0; \frac{1}{2}) \oplus (\frac{1}{2}; 0)$$

AS REP OF $\mathfrak{d}_\ell^\circ = \mathfrak{so}(3) \oplus \mathfrak{su}(2)_R$

$M > |\mathbb{Z}|$: LONG REPS

$$\rho_{hh} \otimes \rho_{hh} \otimes \mathfrak{h}$$

$M = |\mathbb{Z}|$: SHORT REPS: $\mathbb{R}_\alpha^A = 0$,

$$\rho_{hh} \otimes \mathfrak{h}$$

$\mathfrak{h} = \text{ANY FIN. DIML. REP OF } \mathfrak{so}(3) \oplus \mathfrak{su}(2)_R$

PROTECTED SPIN CHARACTER

$$\text{ch}(\rho) = \text{Tr}_\rho X_1^{2J_3} X_2^{2I_3}$$

$$= \begin{cases} (X_1 + X_1^{-1} + X_2 + X_2^{-1})^2 \text{ch}(\mathfrak{h}) & \text{LONG} \\ (X_1 + X_1^{-1} + X_2 + X_2^{-1}) \text{ch}(\mathfrak{h}) & \text{SHORT} \end{cases}$$

\therefore

$$X_1 \frac{\partial}{\partial X_1} \left(\text{Tr}_\rho (X_1^{2J_3} X_2^{2I_3}) \right) \Big|_{X_1 = -X_2 = y}$$

$$= \text{Tr} (2J_3) (-1)^{2J_3} (-y)^{2(J_3 + I_3)}$$

$$= (y - y^{-1}) \text{ch}(\mathfrak{h}) \Big|_{X_1 = -X_2 = y}$$

IS AN INDEX!

MODULI OF VACUA

1. $\mathcal{N}=2$ THEORIES HAVE
MODULI SPACES OF VACUA

2. LOCALLY

$$\left\{ \begin{array}{c} \text{VM} \\ \text{MODULI} \end{array} \right\} \times \left\{ \begin{array}{c} \text{HM} \\ \text{MODULI} \end{array} \right\}$$

↑
THIS TALK : \mathcal{B}

3. LOW ENERGY THEORY HAS
UNBROKEN RANK r

SELF DUAL ABELIAN
GAUGE THEORY

SELF DUAL ABELIAN GAUGE THEORY

1. LOCAL SYSTEM $\Gamma \rightarrow \mathcal{B}$
OF RANK $2r$ SYMPLECTIC
LATTICES

(--- MORE GENERALLY Γ IS
POISSON:

$$0 \rightarrow \Gamma_{\text{flavor}} \rightarrow \Gamma \rightarrow \Gamma_{\text{gauge}} \rightarrow 0$$

2. $V = \Gamma \otimes \mathbb{R}$ HAS COMPATIBLE
POSITIVE COMPLEX STR. \mathbb{I}

- $\mathbb{F} \in \Omega^2(M^4) \otimes V$
- $\mathbb{F} = (* \otimes \mathbb{I}) \mathbb{F}$
- $d\mathbb{F} = 0$.

3. \nexists LAGRANGIAN, BUT
CHOICE OF DUALITY FRAME

(α_I, β^I) FOR $\Gamma \Rightarrow$

• $I \iff \tau_{IJ}$

• $\frac{1}{4\pi} \int \text{Im} \tau_{IJ} F^I * F^J + \text{Re} \tau_{IJ} F^I F^J$

Remark: ALSO IN $D=6$ ABEL. TM

ONCE ONE CHOOSES A LAG.

DECOMPOSITION OF THE SPACE
OF FIELDS \exists AN ACTION
PRINCIPLE (BELOV $\&$ MOORE)

CONSTRAINTS OF SUSY

$$\mathcal{H} = 1\text{-PARTICLE H.S.}$$

$$= \bigoplus_{\gamma \in \Gamma} \mathcal{H}_\gamma$$

RESTRICTION OF \hat{Z} OPERATOR
TO \mathcal{H}_γ IS A SCALAR \Rightarrow

$Z(\gamma)$: CENTRAL CHARGE
FUNCTION

$$N=2 \text{ SUSY} \Rightarrow$$

(1) $Z(\gamma)$ DEPENDS ON VAC.
MODULI BUT ONLY ON VM: $Z(\gamma; u)$

$$(2) Z(\cdot; u) \in \text{Hom}(\Gamma, \mathbb{C})$$

LIES IN A LAGRANGIAN SUBMFLD:

$$Z(\beta^{\mathbb{I}}) = -a^{\mathbb{I}} \quad Z(\alpha_{\mathbb{I}}) = a_{D, \mathbb{I}} = \frac{\partial \mathcal{F}}{\partial a^{\mathbb{I}}}$$

DEF. OF PROTECTED SPIN CHARACTER

$$\mathcal{H}^{\text{BPS}} = \bigoplus_{\gamma \in \Gamma} \mathcal{H}_{\gamma}^{\text{BPS}} \leftarrow \begin{array}{l} \text{FINITE} \\ \text{DIM'L} \end{array}$$

$$(y - \bar{y}^{-1}) \Omega(\gamma; u; y) := \text{Tr}_{\mathcal{H}_{u, \gamma}^{\text{BPS}}} (2J_3)(-1) (-y)^{2J_3}$$

$$\Omega(\gamma; u; y) = \text{ch}(\mathcal{H}_{\gamma}^{\text{BPS}}) \Big|_{x_1 = -x_2 = y}$$

DEF: THE BPS INDEX IS

THE SPECIALIZATION TO $y = -1$:

$$\Omega(\gamma; u) := \Omega(\gamma; u; y) \Big|_{y = -1}$$

IMPORTANT REMARK

IN SUGRA THERE IS IN GENERAL
NO $SU(2)_R$ SYMMETRY.

BUT BPS INDICES CAN
STILL BE DEFINED, BECAUSE

$$\begin{aligned}\Omega(\gamma; u) &= -\frac{1}{2} \text{Tr}_{\mathcal{H}_{\gamma, u}^{\text{BPS}}} (2J_3)^2 (-1)^{2J_3} \\ &= \text{Tr}_{\mathcal{H}_{\gamma, u}} (-1)^{2J_3}\end{aligned}$$



∃ SATISFACTORY PHYSICAL DERIVATION
OF KSWCF FOR BPS INDICES IN BOTH
FIELD THEORY & SUGRA, BUT ONLY FOR
THEORIES WITH UNBROKEN $SU(2)_R$ IS THERE
A PHYSICAL DERIVATION OF THE
MOTIVIC W.C.F.

OPEN PROBLEM: FIND AN
ALGORITHM TO COMPUTE THE
BPS SPECTRUM

- * NO EXAMPLES IN SUGRA
- * \exists FOR "A₁-THEORIES IN CLASS Δ "
- * OTHERWISE UNKNOWN IN F.T.,
EXCEPT FOR SCATTERED EXAMPLES

FIRST STEP: HOW DOES THE
SPECTRUM DEPEND ON MODULI?

- SPIN CHARACTER DEPENDS
ON BOTH VM & HM
- PSC & INDEX ONLY DEPEND
ON VM .

BPS STATES CAN FORM
BPS BOUNDSTATES

FCIV,
SW, ...

BINDING ENERGY

$$\text{B.E.} = |Z(\gamma_1 + \gamma_2)| - |Z(\gamma_1)| - |Z(\gamma_2)| \leq 0$$

DECAYS ONLY ALONG

$$\text{MS}(\gamma_1, \gamma_2) := \{u \mid Z(\gamma_1; u) \parallel Z(\gamma_2; u)\}$$

(B.E. < 0 NEC. NOT SUFF. !)

$\Omega(\gamma; u)$ ONLY JUMPS ACROSS
MS(γ_1, γ_2) WITH $\gamma_1 + \gamma_2 = \gamma$

WCF: FIND $\Delta\Omega$

III. PRIMITIVE & SEMI-PRIMITIVE WCF

(A) M.C. SOLN'S (DENEFF)

DATA: $u = \lim_{\vec{x} \rightarrow \infty} u(\vec{x}) \in \mathcal{B}$

Collection (\vec{x}_j, γ_j) $j=1, \dots, N$ s.t.

$$\forall i \quad \sum_{j \neq i} \frac{\langle \gamma_i, \gamma_j \rangle}{|\vec{x}_{ij}|} = 2 \operatorname{Im} \left(e^{-i\alpha} Z_{\gamma_i}(u) \right)$$

$$e^{i\alpha} = \operatorname{Phase} \left(Z(\sum \gamma_i, u) \right)$$

Solution is a "BPS molecule" of dyonic BPS black holes.

(B) $N=2$: BOUND STATE RADIUS:

$$R_{12} = \frac{1}{2} \langle \gamma_1, \gamma_2 \rangle \frac{|z_1 + z_2|}{\operatorname{Im}(z_1 \bar{z}_2)}$$
$$= \frac{1}{2} \langle \gamma_1, \gamma_2 \rangle \frac{1}{\operatorname{Im}(e^{-i\alpha} z_1)}$$

REMARKS:

(1) DENEFF STABILITY:

$$\langle \gamma_1, \gamma_2 \rangle \operatorname{Im}(z_1 \bar{z}_2) > 0$$

NECESSARY BUT NOT SUFFICIENT

(2.) FOR $Z_2 \rightarrow \infty$

- $|Z_1 + Z_2| - |Z_2| \rightarrow \operatorname{Re}(e^{-i\alpha} Z_1)$
 $= -\operatorname{Re}(Z_1/\gamma)$

FOR $\gamma = -e^{+i\alpha} = -e^{i\alpha_2}$

- 2nd formula is valid provided $\langle \alpha_1, \alpha_2 \rangle$ remains finite

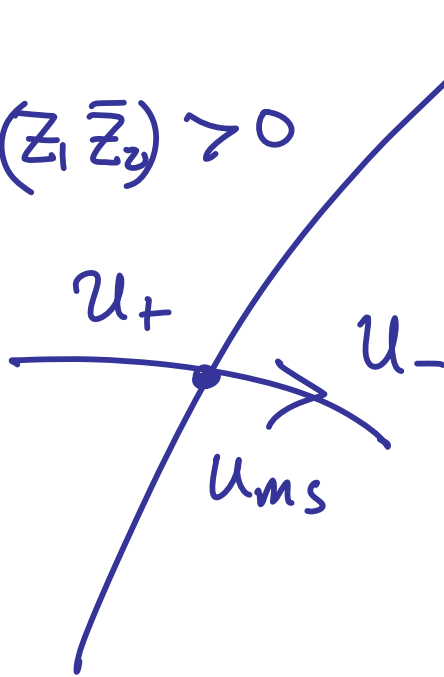
(C.) BASIC WC MECHANISM:

(Denef + Moore)



$$\langle \gamma_1, \gamma_2 \rangle \operatorname{Im}(z_1 \bar{z}_2) > 0$$

$$\langle \gamma_1, \gamma_2 \rangle \operatorname{Im}(z_1 \bar{z}_2) < 0$$



Radius $R_{12} \nearrow \infty \Rightarrow$

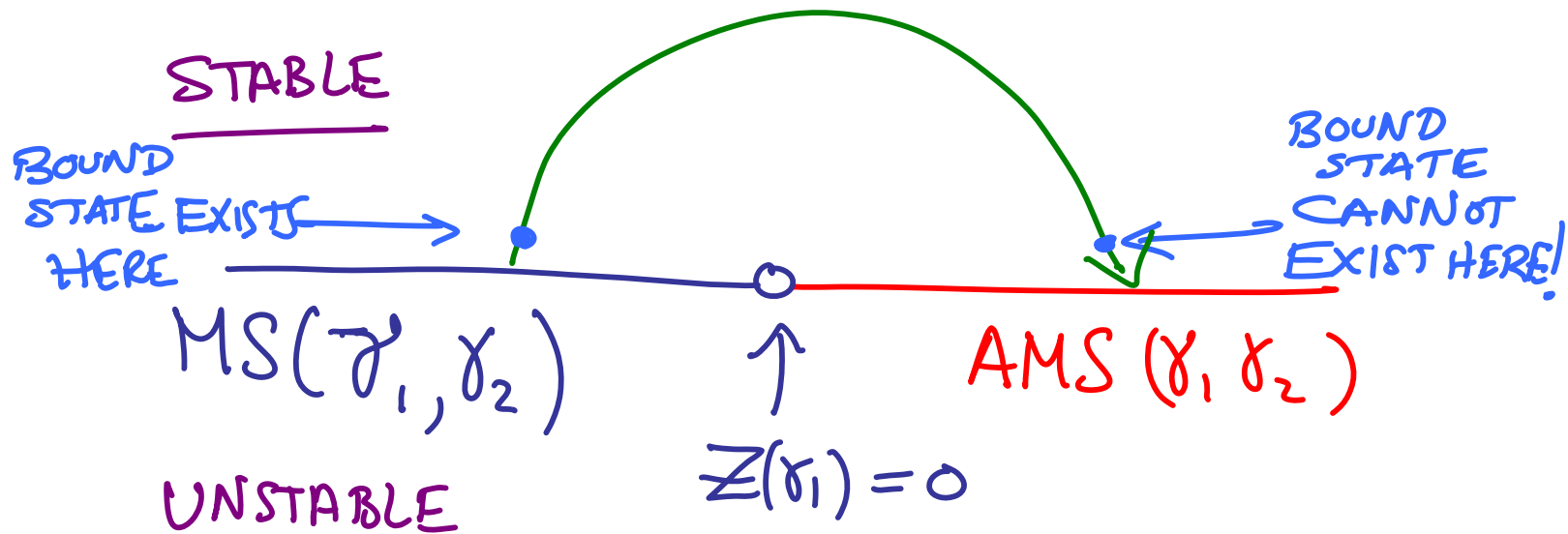
$$\Delta\Omega = \operatorname{Ch}(J_{\gamma_1, \gamma_2}) \Omega_{\bullet}(\gamma_1, u_{ms}, \gamma) \Omega(\gamma_2, u_{ms}, \gamma)$$

γ_1, γ_2 PRIMITIVE

$$2J_{\gamma_1, \gamma_2} + 1 = |\langle \gamma_1, \gamma_2 \rangle|$$

(D.) AMS PARADOX

THIS PICTURE RAISES A PARADOX:



- \exists NO OTHER WALLS OF MS
- γ_1 -PARTICLE IS STABLE IN A NBD \mathcal{U} OF $Z(\gamma_1) = 0$.

RESOLUTION LATER

(E.) HALOS + HALO FOCK SPACES

ONE REASON THE PRIMITIVE WCF IS INCOMPLETE IS THAT

$$MS(\gamma_1, \gamma_2) = MS(N_1 \gamma_1, N_2 \gamma_2)$$

SO WHEN u CROSSES THE WALL LOTS OF OTHER STUFF DECAYS.

DEF: HALO CONFIG. IS A MC BOUNDSTATE WITH

$$\vec{x} = 0, \gamma_c$$

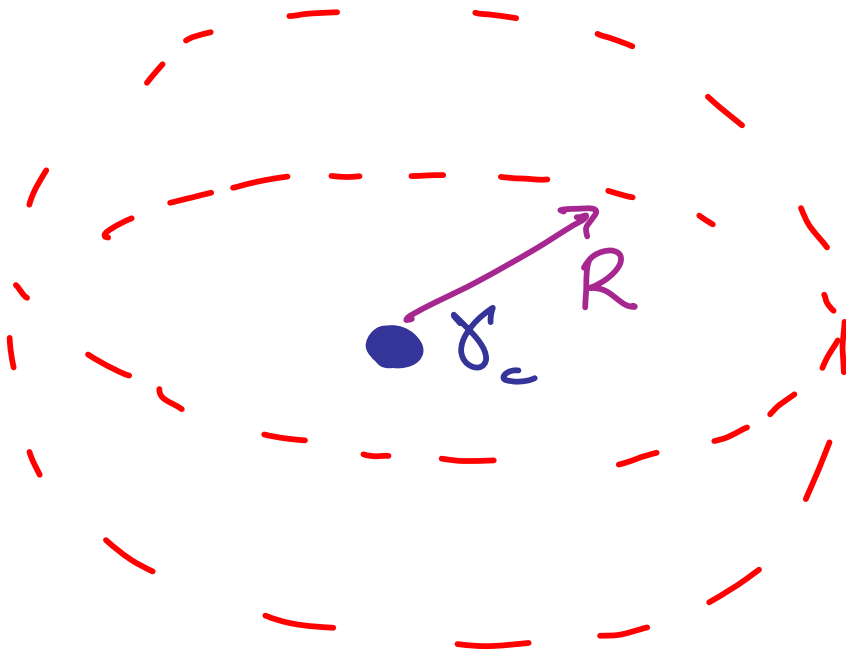
$$\vec{x}_i, \gamma_i = \lambda_i \gamma_h \quad \lambda_i > 0$$

CONSTRAINTS \Rightarrow ALL \vec{x}_i

SIT ON SPHERE:

$$R = \frac{1}{2} \langle \gamma_c, \gamma_h \rangle \frac{1}{\text{Im}(e^{-i\alpha} Z_{\gamma_h})}$$

AND NO OTHER CONSTRAINTS APPLY



HALO PARTICLES ARE MUTUALLY
BPS \Rightarrow QUANTUM STATES
ARE N -PARTICLE FOCK SPACE
STATES.

ONE PARTICLE SPACE OF
QUANTUM STATES: of a halo particle

$$\mathcal{H}_{u, \lambda \gamma_h}^{\text{BPS}} = \mathcal{P}_{hh} \otimes \mathcal{H}_{\lambda \gamma_h}$$

$\rho_{hh} \sim$ COM. MUST BE
IN RADIAL SPIN $\frac{1}{2}$ STATE

\Rightarrow ~~ρ_{hh}~~ BUT ON THE OTHER HAND

THERE ARE LANDAU-LEVEL DEGENERACIES
OF THE QUANTUM STATES WHICH ARE
HENCE DRAWN FROM

$$W_{\lambda\gamma_h} = (J_{\gamma_c, \lambda\gamma_h}) \otimes \mathbb{Z}_{\lambda\gamma_h}$$

$$\mathbb{Z}_2 - \text{GRADING} \quad -(-1)^{2J_3^h \lambda\gamma_h}$$

HM's \sim Halo fermions

VM's \sim Halo bosons

The quantum states associated with halo configurations with core γ_c is:

$$\bigoplus_{N \geq 0} q^N \mathcal{H}_{\gamma_c + N\gamma}^{\text{HALO}} =$$

$$\mathcal{H}_{\gamma_c} \otimes_{n=1}^{\infty} \mathcal{F}[q^n W_{n\gamma_n}]$$

$\mathcal{F}[W] = \mathbb{Z}_2$ -GRADED FOCK SPACE
 GENERATED BY \mathbb{Z}_2 -GRADED
 V.S. W

(F.) SEMI-PRIMITIVE WCF.

NOW WRITE A WCF FOR THE
HALO CONTRIBUTION TO THE
INDEX:

Introduce group algebra for Γ

$$X_{\gamma_1} X_{\gamma_2} = X_{\gamma_1 + \gamma_2}$$

$$G_{\gamma_c}^{\text{Halo}} = \sum_{N=0}^{\infty} \Omega_{\gamma_c}^{\text{Halo}}(N\gamma_h) X_{\gamma_c + N\gamma_h}$$

UNSTABLE

$$G_{\gamma_c}^{\text{Halo}} = X_{\gamma_c}$$

STABLE :

$$G_{\gamma_c}^{\text{Halo}} = \prod_{n=1}^{\infty} \left(1 - (-1)^{\langle n\gamma_h, \gamma_c \rangle} X_{\gamma_h} \right) \cdot X_{\gamma_c}$$

$\langle n\gamma_h, \gamma_c \rangle \mid \Omega(n\gamma_h \text{ illus})$

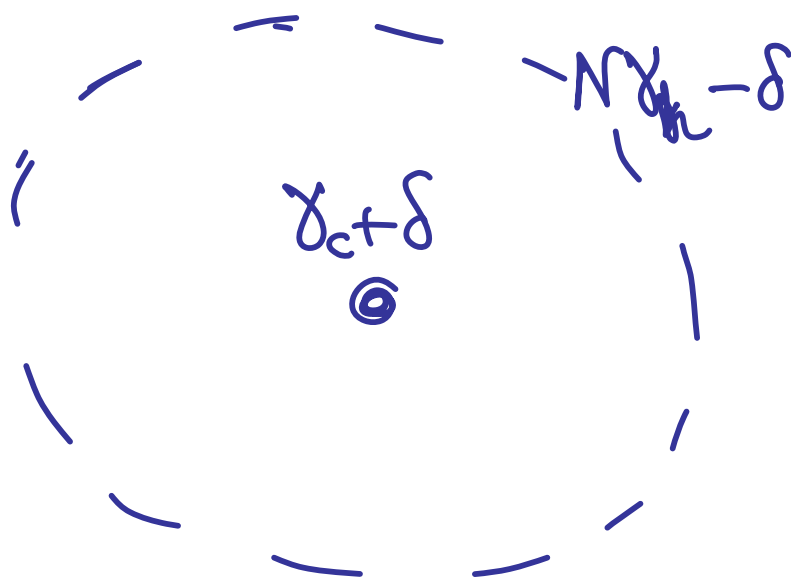
"SEMI-PRIMITIVE WCF"

IV. PHYSICAL DERIVATION OF KSWCF

(A.) MAIN PROBLEM W/ HALO
PICTURE IS THAT

$$\Omega(\gamma_c + N\gamma) \neq \Omega(\gamma_c) \Omega_{\gamma_c}^{\text{Halo}}(N\gamma)$$

⌋ Mixing:



WITH OTHER MORE COMPLICATED
BOUND STATES OF TOTAL CHARGE

$$\gamma_c + N\gamma$$

(B.) BPS GALAXIES

NEW IDEA: SUPPRESS MIXING BY
TAKING LARGE CORE CHARGE

CHOOSE $U(1)$ SUBGROUP OF
RANK r ABELIAN GAUGE GROUP
(el., mag.) CHARGES γ_0', γ_0 ,

$$\langle \gamma_0, \gamma_0' \rangle = 1$$

TAKE: $\gamma_c = \Lambda^2 \gamma_0 + \Lambda \gamma_0' + \delta$

$$\delta \in \Gamma_0^\perp := \{ \gamma \mid \langle \gamma, \gamma_0 \rangle = \langle \gamma, \gamma_0' \rangle = 0 \}$$

WE WILL TAKE $\Lambda \rightarrow \infty$

WE DEFINE A "BPS GALAXY"

TO BE A M.C. SOLUTION
(OR ITS QUANTUM ANALOG)

WHERE ONE CENTER $\vec{x} = 0$ HAS
CHARGE γ_c AND ALL OTHER
CENTERS \vec{x}_j HAVE CHARGES
 $\gamma_j \in \Gamma_0^\perp$.

KEY CLAIM: LET

$\mathcal{M}(\gamma_c) = \left\{ \begin{array}{l} \text{ENSEMBLE OF ALL} \\ \text{BPS GALAXIES} \end{array} \right\}$

FOR $\Lambda \rightarrow \infty$ THERE IS NO
QUANTUM MIXING BETWEEN
ENSEMBLES WITH $\gamma_c \neq \gamma'_c$

JUSTIFICATION

(1.) ENTROPIC SUPPRESSION:

BIG BLACK HOLES CAN'T FRAGMENT

cf. Maldacena-Michelson-Strominger

(2.) DISTANCE SUPPRESSION:

TUNNELING OVER LARGE DISTANCES
IS EXPONENTIALLY SUPPRESSED

(C.) FRAMED BPS INDEX

$$h_{\gamma_c}(\gamma_{\text{orb}}; u) := h_{\gamma_c + \gamma_{\text{orb}}} / h_{\gamma_c}$$

$$\overline{\overline{\Sigma}}_{\gamma_c}(\gamma_{\text{orb}}; u) := \lim_{\Lambda \rightarrow \infty} \text{Tr}_{h_{\gamma_c}(\gamma_{\text{orb}}; u)} (-1)^{2J_3}$$

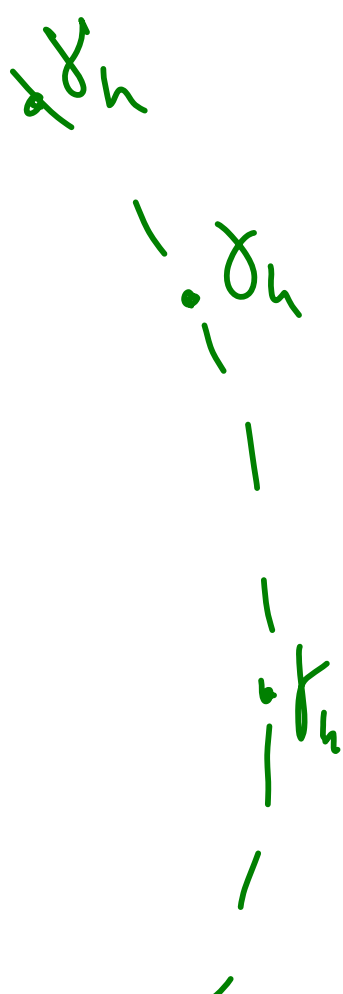
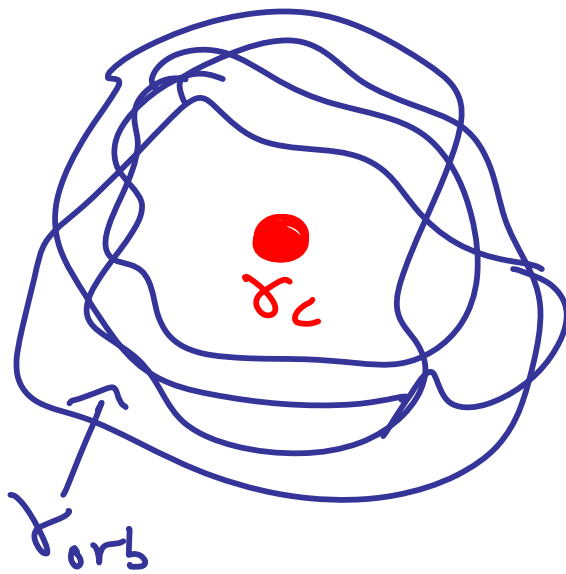
Claim These are well-defined and piecewise constant.

(D.) Framed BPS Wall-Crossing

BPS Galaxies are complicated
But Wall-crossing is simple:

$$Z(\gamma_{hiu}) \parallel Z(\gamma_{tiu})$$

\Rightarrow haloes $\gamma_t := \gamma_c + \gamma_{orb}$
galaxy: — —



HAPPENS WHEN

$$Z(\gamma_h) \parallel Z(\gamma_c + \gamma_{orb})$$

$\Lambda \rightarrow \infty$:

$$W_{\gamma_h} = \{u \mid Z(\gamma_h) \parallel Z(\gamma_o)\}$$

N.B. DEPENDENCE ON γ_{orb} DROPPED OUT!

NOW STUDY:

$$G_{\gamma_c} = \sum_{\gamma_{orb} \in \mathbb{T}_o^\perp} \overline{\Omega}_{\gamma_c}(\gamma_{orb}; u) \times_{\gamma_t} \cdot \begin{pmatrix} -1^2 & -1 \\ X_{\gamma_o} & X_{\gamma_o'} \end{pmatrix}$$

HALOS \Rightarrow

$$X_{\gamma_t} \rightarrow \prod_n \left(1 - (-1)^{\langle n \gamma_h, \gamma_t \rangle} X_{\gamma_h}^n \right)^{|\langle n \gamma_h, \gamma_t \rangle| \Omega(\gamma_h; u)} \times_{\gamma_t}$$

Note $\gamma_h \in \Gamma_\epsilon^\perp$

$$\langle \gamma_h, \gamma_t \rangle = \langle \gamma_h, \delta + \gamma_{\text{orb}} \rangle$$

IS WELL-DEFINED FOR $\Lambda \rightarrow \infty$.

BUT! THE HALO FACTOR
DEPENDS ON γ_{orb} .

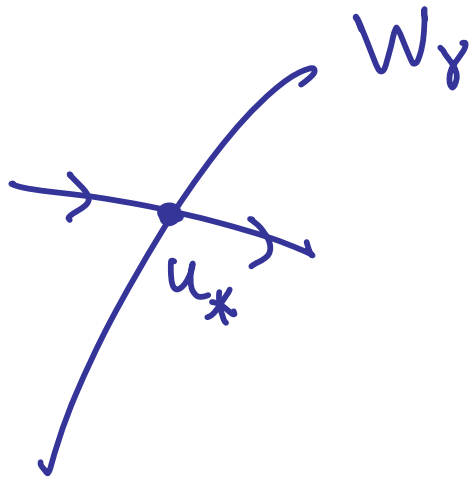
\Rightarrow INTRODUCE

$$D_\gamma X_{\gamma'} := \langle \gamma, \gamma' \rangle X_{\gamma'}$$

$$K_\gamma := \left(1 - (-1)^{D_\gamma} X_\gamma \right)^{D_\gamma}$$

Then across a BPS Wall

$$G_{\gamma_c} \rightarrow U_{\gamma}(u_*) \cdot G_{\gamma_c}$$

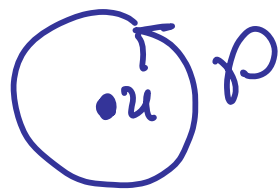


$$U_{\gamma}(u_*) = \prod_{n=1}^{\infty} K_{n\gamma} \Omega(n\gamma; u_*)$$

IS A DIFFL. OPERATOR.

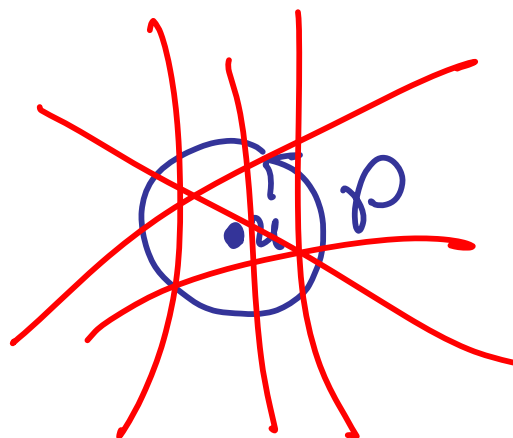
(E.) DERIVING KSWCF

CONSIDER A NONSINGULAR
POINT IN MODULI SPACE AND
A SMALL PATH \mathcal{P} NEARBY:



$G_{\gamma_c}(u)$ IS SINGLE-VALUED ON \mathcal{P}

BUT THERE WILL TYPICALLY
BE MANY WALLS W_{γ}



So

$$G_{\gamma_c}(u) \rightarrow \prod U_{\gamma_i}(u_i) \cdot G_{\gamma_c}(u)$$

SINCE δ WAS ARBITRARY

$G_{\gamma_c}(u)$ IS ARBITRARY:

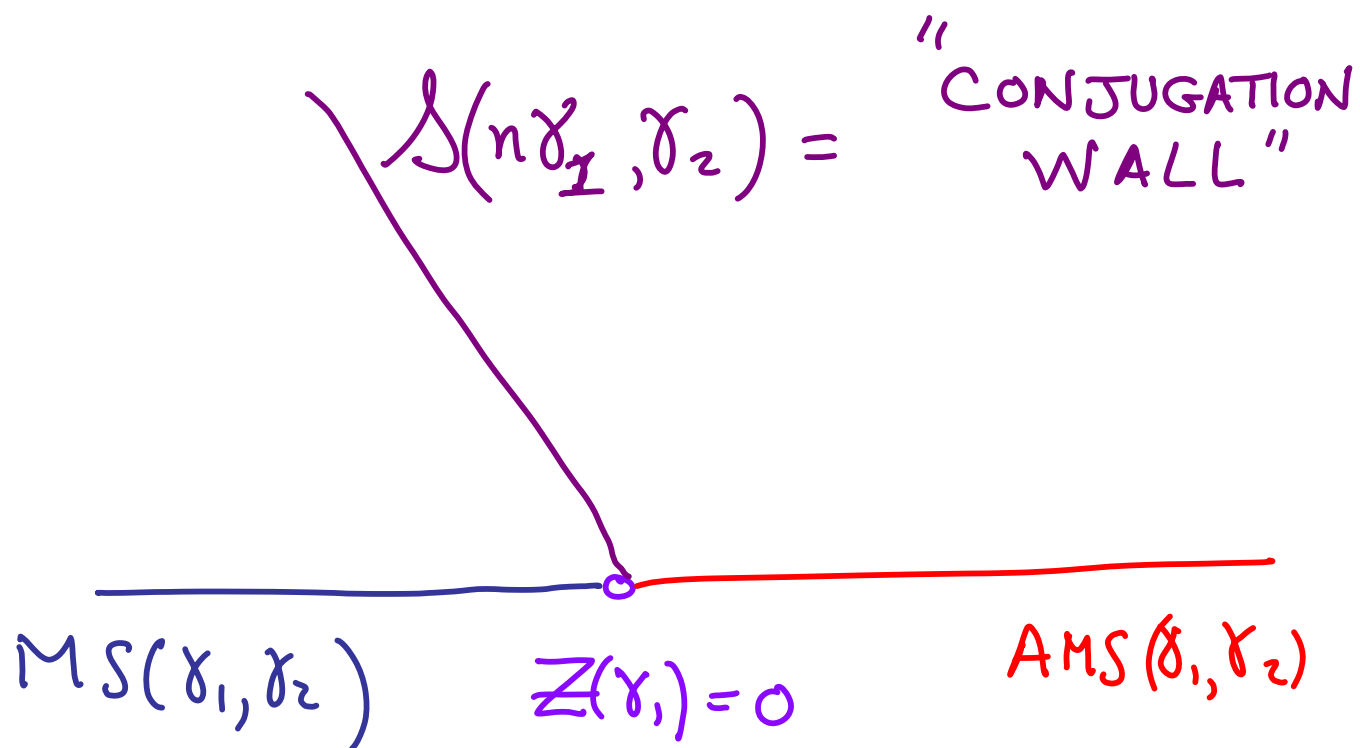
$$\prod_i U_{\gamma_i}(u_i) = 1$$

OPERATOR EQUATION CONSTRAINS
THE BPS DEGEN'S $\Omega(\gamma_i)$!

IT IS EASY TO SHOW
THIS IMPLIES KSWCF.

(F.) AMS PARADOX

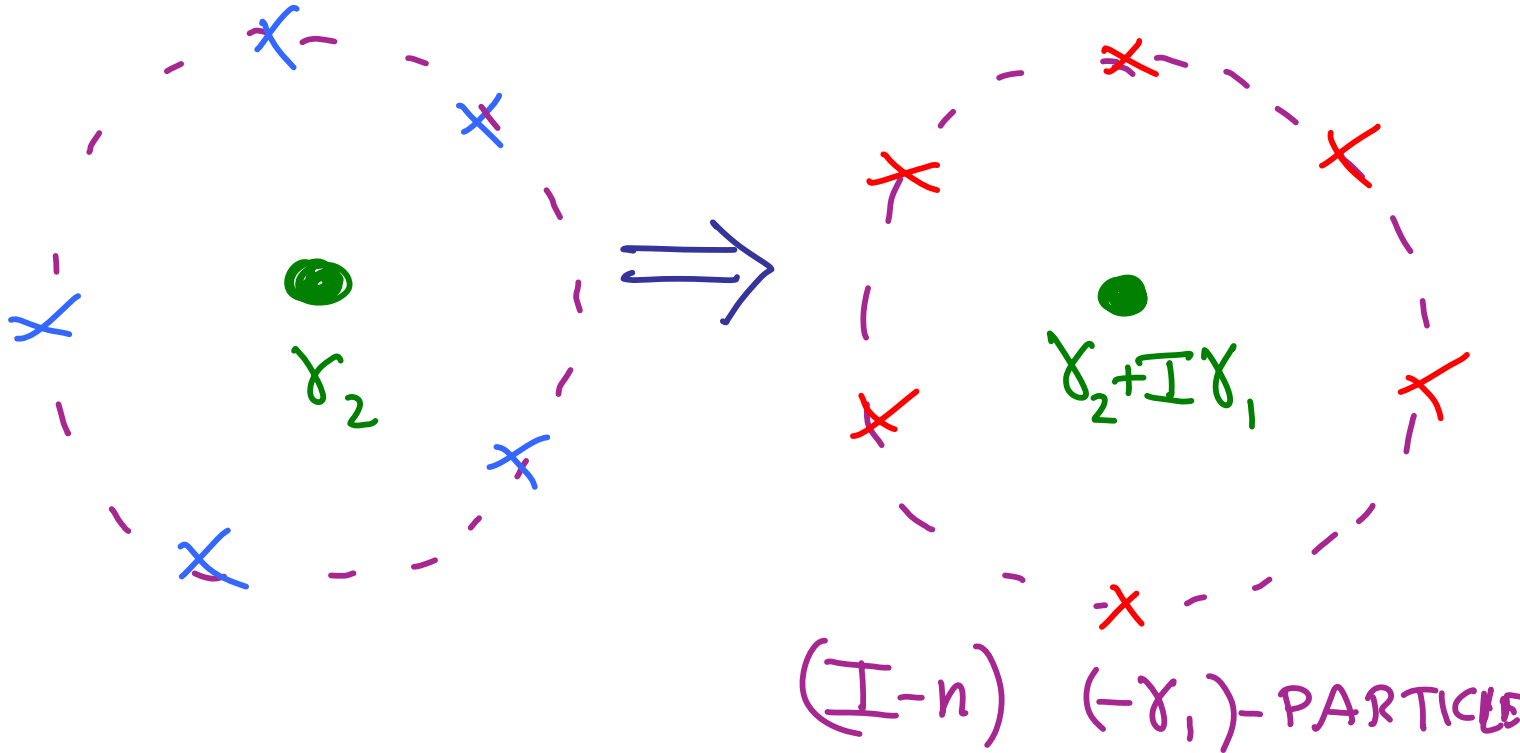
\exists NEW WALL WHERE HALO
DESCRIPTION CHANGES



$S(n\gamma_1, \gamma_2)$ IS DEFINED IN
TERMS OF ATTRACTOR FLOW TREES:

$\{u \mid \text{FLOW OF } n\gamma_1 + \gamma_2 \text{ CRASHES ON } Z(\gamma_1) = 0\}$

n γ_1 -PARTICLES



$\mathbb{I} = \text{MONODROMY:}$

$$\gamma_2 \rightarrow \gamma_2 + \mathbb{I} \gamma_1$$

$$\mathbb{I} = |\langle \gamma_1, \gamma_2 \rangle| \sum_{k=1}^{\infty} k^2 \Omega(k\gamma) > 0.$$

TWO CONSEQUENCES

① IF THE ONLY MASSLESS POPULATED CHARGES AT $\mathcal{Z}(\gamma)$ ARE \parallel TO γ THEN:

$$P(q) = \prod_k (1 - q^k)^{kS_2(k\gamma)} \in \mathbb{Z}[q]$$

NOTA BENE: MASSLESS VM'S

CONTRIBUTE $(1 - q^k)^{-2k}$!

② WHEN THERE ARE MASSLESS VM'S THERE ARE ALSO MASSLESS MONOPOLES !

V LINE DEFECTS

(A) NOW WE FOCUS ON FIELD THEORY

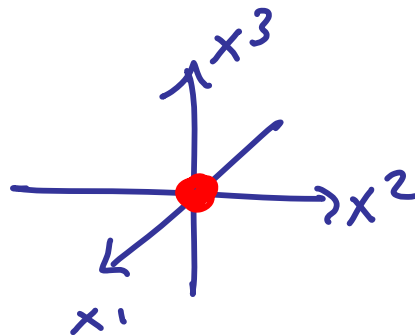
BASED ON SCFML FXD PT \mathcal{L}_*

LINE DEFECT: NBD IS CONFORMAL
TO $AdS_2 \times S^2$

DEFINE IT TO BE A CONFORMAL
BDRY CONDITION OF \mathcal{L}_* [KAPUSTIN]

CONSIDER A LINE DEFECT

$$\{\vec{X} = 0\} \times \mathbb{R}$$



DEF: WE SAY L IS OF TYPE \mathcal{J} IF IT PRESERVES THE SUBALGEBRA FIXED BY \mathcal{J}

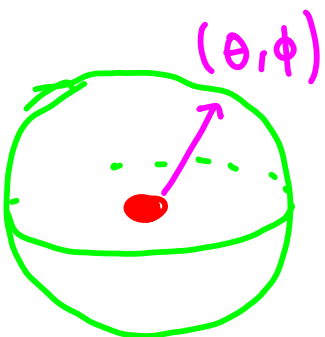
$$R_\alpha^A \sim Q_\alpha^A + \mathcal{J} \sigma_{\alpha\beta}^0 \bar{Q}^{\beta A}$$

EXAMPLES

1.) WILSON

$$L_{\mathcal{J}} \sim P \exp \int_{\mathbb{R} \times \vec{0}} \left(\frac{\varphi}{2\mathcal{J}} - iA - \frac{\mathcal{J}}{2} \bar{\varphi} \right)$$

2.) 't HOOFT



$$F \sim p \otimes \sin\theta d\theta d\phi \quad p \in \mathfrak{k}$$

$$\varphi/\mathcal{J} \sim \frac{p}{r} + \varphi_\infty/\mathcal{J}$$

(B.) FRAMED BPS STATES

HILBERT SPACE IS MODIFIED

$$\mathcal{H}_L = \bigoplus_{\gamma \in \Gamma_L} \mathcal{H}_{L,\gamma}$$

(IN GENERAL $\Gamma_L = \Gamma - \underline{\text{TORSOR}}$)

THINK OF $\mathcal{H}_{L,\gamma}$ AS SECTOR
CREATED BY ∞ -LY HEAVY DYON
OF CHARGE $\gamma \Rightarrow$

$$E \geq -\text{Re}(Z_\gamma/s)$$

DEF: FRAMED BPS STATES

SATURATE THIS BOUND

DEF: THE FRAMED

PROTECTED SPIN CHARACTER

$$\underline{\bar{\Omega}} := \text{Tr}_{\mathcal{H}_{L_S, \gamma, u}^{\text{BPS}}} (-1)^{2J_3} (-y)^{2g_3}$$

$$g_3 = J_3 + I_3$$

$\underline{\bar{\Omega}} \in \mathbb{Z}[y, y^{-1}]$ AND

$$\underline{\bar{\Omega}}(L, \gamma; y; \mathcal{S}; u)$$

IS PIECEWISE CONSTANT

IN \mathcal{S}, u

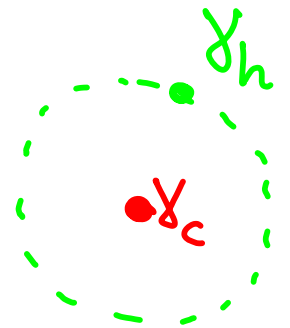
(C.) FRAMED BPS WALL-CROSSING

NEAR BPS WALLS:

$$W(\gamma_h) = \left\{ (u, \mathcal{S}) \mid \Re z_{\gamma_h} / \mathcal{S} < 0 \right\} \\ \subset \mathcal{B} \times \mathbb{C}^*$$

SOME OF THE FRAMED BPS STATES ARE DESCRIBED BY HALOS:

$$r_{\text{HALO}} = \frac{\langle \gamma_h, \gamma_c \rangle}{2 \operatorname{Im} (z_{\gamma_h} / \mathcal{S})}$$



(FOLLOWS FROM LIMIT OF BOUNDSTATE RADIUS IN THE LIMIT $z_2 \rightarrow \infty$.)

CROSS $W(\gamma_h)$: CREATE/ANN. HALO
FOCK SPACES

THE BPS WALLS $W(\gamma_h)$

DIVIDE $\mathcal{B} \times \mathbb{C}^*$ INTO CHAMBERS

GENERATING FUNCTION:

$$F(L, c) := \sum_{\gamma} \underline{\underline{\Omega}}(L; \gamma; y; c) X_{\gamma}$$

$$X_{\gamma_1} X_{\gamma_2} = y^{\langle \gamma_1, \gamma_2 \rangle} X_{\gamma_1 + \gamma_2}$$

$\text{Im}(z_{\gamma_h}/s) > 0$	$W(\gamma_h)$	$\text{Im}(z_{\gamma_h}/s) < 0$
C_+		C_-
$F(L, C_+)$		$F(L, C_-)$

FOCK SPACE COMBINATORICS
ARE SUMMARIZED BY

$$F(L, c_+) = S_{\gamma_h} F(L, c_-) S_{\gamma_h}^{-1}$$

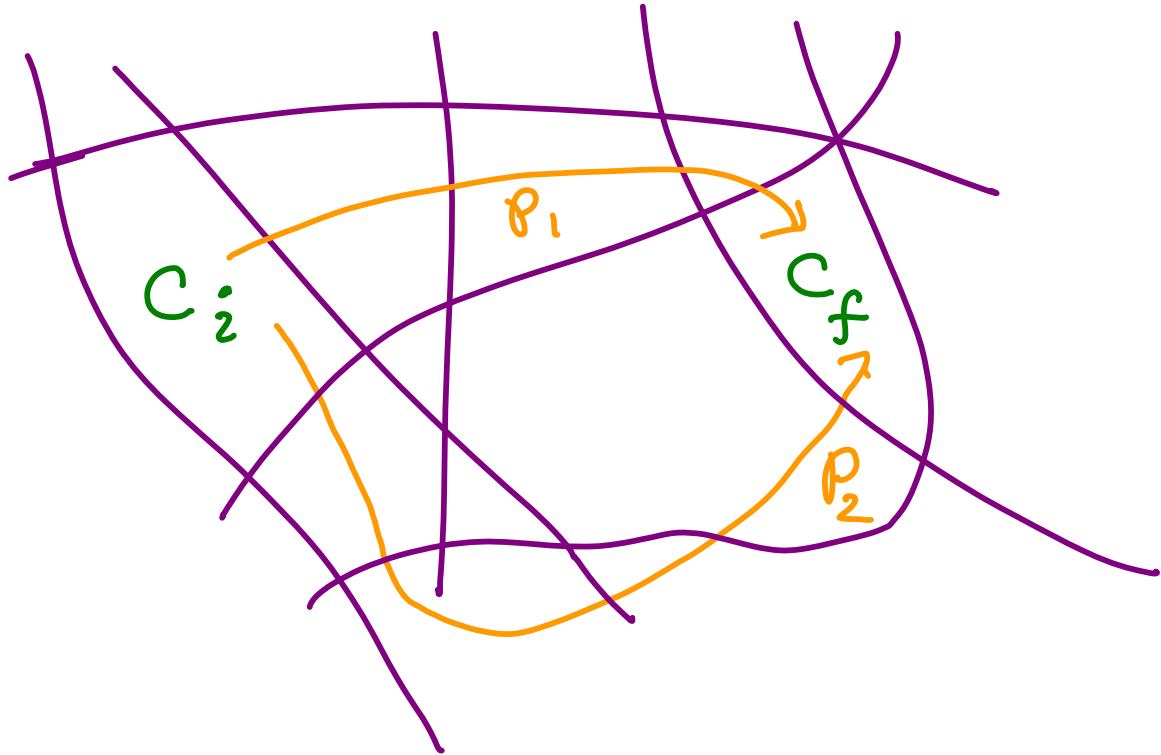
$$S_{\gamma_h} = \prod_{m=-M_h}^{M_h} \underline{\Phi} \left((-y)^m X_{\gamma_h} \right)^{a_{m, \gamma_h}}$$

$$\underline{\Phi}(X) = \prod_{k=1}^{\infty} (1 + y^{2k-1} X) = \text{QUANTUM DILOG}$$

$$\text{Tr}_{\gamma_h} (-1)^{2J_3} (-y)^{2J_3} = \sum_{-M_h}^{M_h} a_{m, \gamma_h} y^m$$

PSC OF HALO PARTICLES

MOTIVIC KSWCF



$$\begin{aligned} F(L, C_f) &= S(\rho_1) F(L, C_i) S(\rho_1)^{-1} \\ &= S(\rho_2) F(L, C_i) S(\rho_2)^{-1} \end{aligned}$$

MOTIVIC
KSWCF: $S(\rho_1) = S(\rho_2)$

RELATED DISCUSSIONS:

DIMOFTE $\stackrel{!}{\approx}$ GUKOV

CECOTTI $\stackrel{!}{\approx}$ VAFA

DIMOFTE, GUKOV $\stackrel{!}{\approx}$ SOIBELMAN

REMARK : $y = -1$

$$\hat{X}_{\gamma_1} \hat{X}_{\gamma_2} = (-1)^{\langle \gamma_1, \gamma_2 \rangle} \hat{X}_{\gamma_1 + \gamma_2}$$

\hat{X}_{γ} : COMMUTATIVE VARIABLES !

$$\begin{aligned} F(L, c_+) &= \sum \overline{\Omega}(L, \gamma) \hat{X}_{\gamma} \\ &= K_{\gamma_h}^{-\Omega(\gamma_h)} F(L, c_-) \end{aligned}$$

$$K_{\gamma_h}(\hat{X}_{\gamma}) := (\hat{X}_{\gamma_h})^{\langle \gamma, \gamma_h \rangle} \hat{X}_{\gamma}$$

(D.) STRONG POSITIVITY

SURPRISING FACT: IN THE
 A_1 -FIELD THEORIES

PSC = SPIN CHARACTER

(NONTRIVIAL STATEMENT
ABOUT TOPOLOGY OF MONOPOLE
MODULI SPACES.)

\Rightarrow PURELY ALGEBRAIC CONSTRUCTION
OF THE ALGEBRA OF LINE DEFECTS

DEF: A STRONGLY POSITIVE
FORMAL LINE OPERATOR IS A
COLLECTION $F(c)$ $c \in \text{CHAMBERS}$

$$\bullet F(c) = \sum_y P_y^c X_y \quad \underline{\text{FINITE SUM}}$$

$P_y^c =$ TRUE SPIN CHARACTER

$$= \sum_{n \geq 0} a_n \frac{y^n - \bar{y}^n}{y - y^{-1}}$$

$$a_n \in \mathbb{Z}_+$$

$$\bullet F(c^+) = S_{\gamma_n} F(c^-) S_{\gamma_n}^{-1}$$

CONJECTURE: FORMAL LINE OPERATORS
COINCIDE WITH THE TRUE LINE OPS

x

- CLUSTER ALGEBRAS
- CLUSTER VARIETIES
- TROPICAL LABELS



VI. HK GEOMETRY $\frac{1}{\epsilon}$

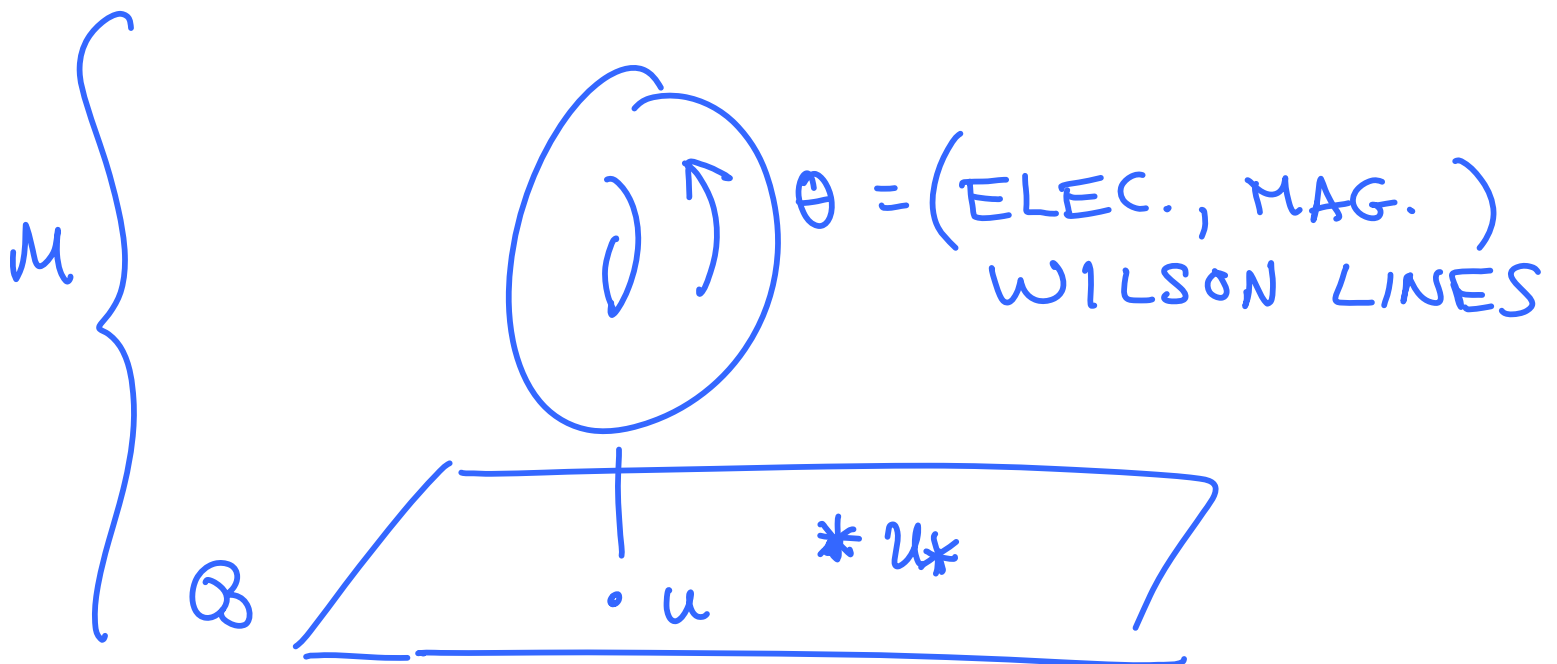
COMPACTIFICATION TO 3DIM'S

(A.) PUT THEORY ON $\mathbb{R}^3 \times S^1_R$

LOW ENERGY EFFECTIVE THEORY:

σ -MODEL : $\mathbb{R}^3 \rightarrow \mathcal{M}$

\mathcal{M} IS A TORUS FIBRATION



SINGULAR ON DISCRIMINANT LOCUS
WHERE POPULATED CHARGES γ ARE
MASSLESS $Z_\gamma \rightarrow 0$.

(B.) "DARBOUX COORD'S"

- SUSY \Rightarrow \mathcal{M} IS HK

- \Rightarrow FAMILY OF CPLX STRUCTURES $\mathcal{M}^S \quad S \in \mathbb{P}^1$

- WRAP L_S ON $S^1 \Rightarrow$

LOCAL OPERATOR IN 3D THEORY

$\langle L_S \rangle$ IS HOLOMORPHIC ON \mathcal{M}^S

- ON THE OTHER HAND

$$\langle L_S \rangle = \text{Tr}_{\mathcal{H}_{L_S}} (-1)^{2J_3} e^{-2\pi R H} e^{i\theta \cdot Q} \sigma(Q)$$

$$\langle L_S \rangle = \text{Tr}_{\mathcal{H}_{L_S}} (-1)^{2J_3} e^{-2\pi R H} e^{i\theta \cdot Q} \sigma(Q)$$

Q = CHARGE OPERATOR

$$\theta \in \text{Hom}(\mathbb{T}, \mathbb{R}/\mathbb{Z}) = \text{B.C.} @ \infty$$

ARE THE WILSON LINES

$\sigma(Q)$: QRIF: SELF-DUALITY!

POSIT AN EXPANSION

$$\langle L_S \rangle = \sum_{\gamma} \overline{\Omega}(L_S, \gamma) \mathcal{Y}_{\gamma}$$

• UV DETAILS OF WHICH LINE OPERATOR CAPTURED BY $\overline{\Omega}(L_S, \gamma)$

• IR EFFECTIVE LINE OPERATOR CREATED BY ∞ -ly HEAVY DYON OF CHARGE $\gamma = \mathcal{Y}_{\gamma} =$ "DARBOUX COORD'S"

(C.) PROPERTIES OF y_γ

(1.) $y_\gamma(\cdot, s)$ HOLOMORPHIC ON M^S

(2.) HOLOMORPHIC IN $s \in \mathbb{C}^*$

(3.) $R \rightarrow \infty$: TRACE ASYMPTOTES
 $\overline{10}$

$$\sum_\gamma \overline{\Omega}(L_{s,\gamma}) \underbrace{e^{\pi R \frac{Z_\gamma}{s} + i \hat{\Theta} \cdot \gamma + \pi R s \overline{Z}_\gamma}}_{y_\gamma^{sf}}$$

SO:

$$y_\gamma \underset{R \rightarrow \infty}{\sim} y_\gamma^{sf}$$

LIKEWISE:

$$y_\gamma \underset{s \rightarrow 0, \infty}{\sim} y_\gamma^{sf}$$

(4.) REALITY CONDITION

$$\overline{y_\gamma(\mathcal{S})} = y_{-\gamma}(-1/\bar{\mathcal{S}})$$

(5.) RING STRUCTURE ON
LINE OP'S \Rightarrow

$$y_{\gamma_1} y_{\gamma_2} = (-1)^{\langle \gamma_1, \gamma_2 \rangle} y_{\gamma_1 + \gamma_2}$$

(6.) $\langle L_{\mathcal{S}} \rangle$ WELL-DEFINED

IN UV $\frac{1}{i}$ UNDERGOES NO

PHASE TRANSITION \Rightarrow

SMOOTH ON $\mathcal{M} \times \mathbb{C}^* \Rightarrow$

SMOOTH ACROSS BPS WALLS $W(\gamma_h)$

\Rightarrow

$$F(L_S) = \sum_{\gamma} \underline{\bar{\Omega}}(L_S, \gamma) \hat{x}_{\gamma}$$

$$\rightarrow K_{\gamma_h}^{-\Omega(\gamma_h)} F(L_S)$$

BUT

$$\langle L_S \rangle = \sum_{\gamma} \underline{\bar{\Omega}}(L_S, \gamma) \psi_{\gamma}$$

$$\rightarrow \langle L_S \rangle$$

THEREFORE,

$$\psi_{\gamma} \rightarrow K_{\gamma}^{\Omega(\gamma_h)} \psi_{\gamma}$$

THESE PROPERTIES

DETERMINE THE ψ_{γ}

UNIQUELY.

(D.) CONSTRUCTING THE \mathcal{Y}_x

CONSTRUCT THESE FUNCTIONS
VIA AN INTEGRAL EQUATION:

$$\log \mathcal{Y}_x = \log \mathcal{Y}_x^{sf}$$

$$+ \sum_{x'} \langle x, x' \rangle \Omega(x') \mathbb{K}_{x, * } \log(1 - \mathcal{Y}_{x'})$$

$$\mathbb{K}_{x, * } f = \int_{\mathcal{L}_{x'}} [ds'] f(s')$$

$$\mathcal{L}_{x'} = \{ s \mid \Re x'/s < 0 \}$$

$$[ds'] = \frac{-1}{4\pi i} \frac{ds'}{s'} \frac{s'+s}{s'-s}$$

FORMALLY IDENTICAL TO ZAMOLODCHIKOV'S
TBA

(E.) CONSTRUCTING HK METRICS

USING STRUCTURE OF TORUS
FIBRATION, AND THE
ALGEBRA OF THE Y_γ ,
WE CONSTRUCT THE
HOLOMORPHIC SYMPLECTIC FORM

$$\begin{aligned}\overline{\omega}_S &= \frac{-i}{2J} \omega_+ + \omega_3 - \frac{i}{2} J \omega_- \\ &= \frac{1}{8\pi^2 R} C_{ij} d \log Y_{\gamma^i} \wedge d \log Y_{\gamma^j} \\ C_{ij} &= \text{inverse to } C^{ij} = \langle \gamma^i, \gamma^j \rangle\end{aligned}$$

RMK: QUANTUM DEFORMATION

$$F(L_S, C) = \sum_{\gamma} \bar{\Omega}(L_S, \gamma, C; y) X_{\gamma}$$

DEFINES A NATURAL
QUANTUM DEFORMATION
OF FUNCTIONS ON \mathcal{M}^S

$$\{y_{\gamma}, y_{\gamma'}\}_{\omega^S} = \langle \gamma, \gamma' \rangle y_{\gamma} y_{\gamma'}$$

THIS GENERALIZES A CONSTRUCTION
OF FOCK-GONCHAROV AND TESCHNER

VII. M5-BRANES & HITCHIN SYSTEMS

(A.) 6D (2,0) THEORIES & CLASS \mathcal{A}

CONSIDER NONABELIAN \mathfrak{g} -THEORY
ON RIEMANN SURFACE C WITH PUNCTURES

BREAK $SO(5)_{\mathbb{R}} \rightarrow SO(2) \oplus SO(3)$

AND PARTIALLY TWIST TO PRODUCE

$d=4, \mathcal{N}=2$ THEORY ON $\mathbb{R}^{1,3}$

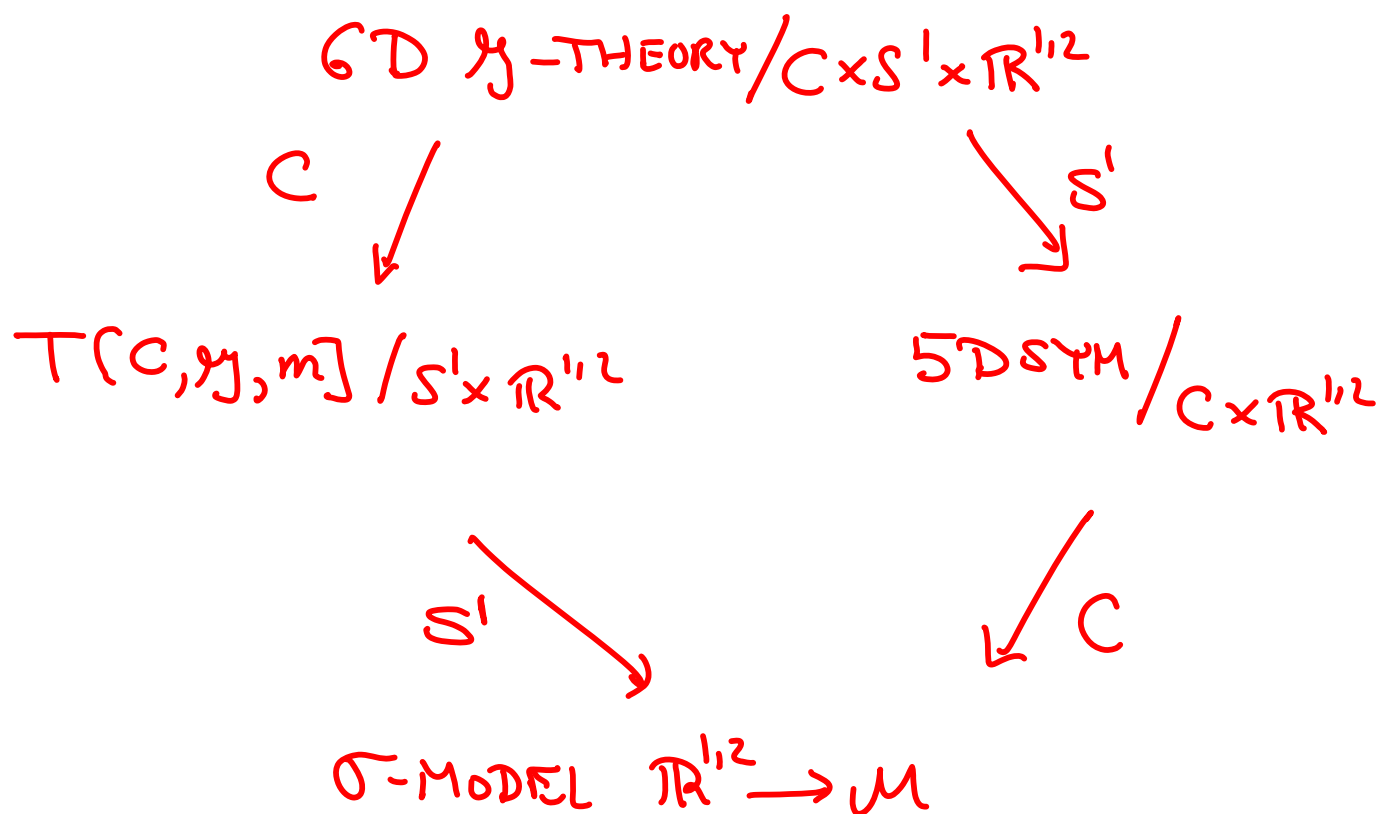
BOUNDARY COND'S AT PUNCTURES
SPECIFIED BY BEHAVIOR OF BPS
OPERATORS \Rightarrow

$T[C, \mathfrak{g}, m]$

THIS IS THE CLASS \mathcal{A} .

(B.) CLASS \mathcal{A} CONTAINS ALL THE FAMILIAR QUIVER THEORIES, AND GENERALIZED QUIVER THEORIES OF GAIOTTO.

WHEN COMPACTIFYING ON S^1 WE CAN IDENTIFY \mathcal{M} WITH MODULI SPACE OF HITCHIN SYSTEMS ON C :



(C.) SW MODULI = HITCHIN MODULI

- REDUCTION OF 5D $\mathcal{N}=2$ SYM ON $C \Rightarrow$ BPS EQS = HITCHIN EQS.

$$F + R^2 [\varphi, \bar{\varphi}] = 0$$

$$\bar{\partial}_A \varphi = 0$$

- SPECTRAL CURVE $\Sigma \subset T^*C$
 $\Sigma \rightarrow C$

= SEIBERG-WITTEN CURVE

- $d\lambda = dq \wedge dp$ ON T^*C
 $\lambda =$ SW DIFFFL

$$\mathcal{M}_{\text{HITCHIN}}^S \cong \mathcal{M}_{\text{FLAT}}(\mathfrak{g}_{\mathbb{C}}\text{-CONN})$$

$$A = \frac{R}{S} \varphi + A + RS \bar{\varphi}$$

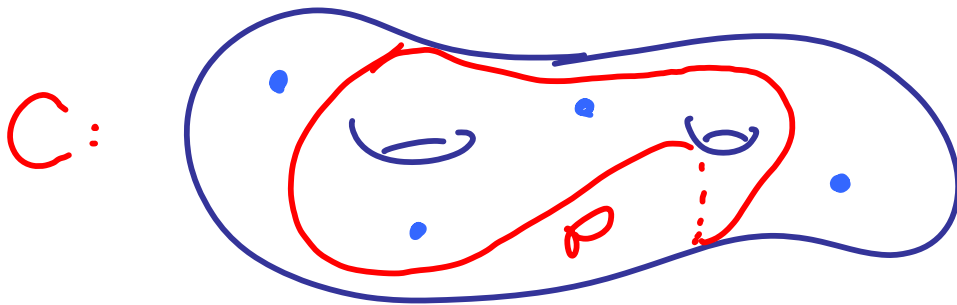
FOR $S \in \mathbb{C}^*$

(C.) A NATURAL SET OF LINE DEFECTS

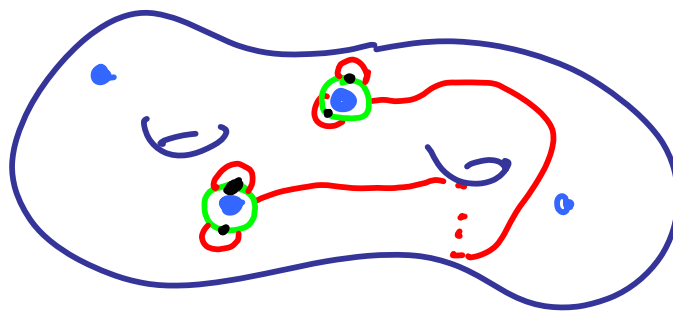
6D $\mathcal{N} = 1$ -THEORY HAS SUSY
SURFACE OP'S $\mathcal{S}(\mathcal{R}, \Sigma)$

$\mathcal{R} \sim \mathcal{N} = 1$ -REP, $\Sigma = \text{SURFACE}$

CHOOSE CLOSED CURVE $\mathcal{P} \subset C$

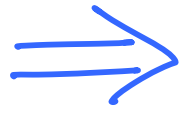


MORE GENERALLY $\mathcal{P} = \text{LAMINATION}$



CHOOSE $\Sigma = S^1_R \times \mathcal{P}$

6D SURFACE DEFECT $\mathcal{S}(\mathcal{R}, \Sigma)$



4D LINE DEFECT $L_S(\mathcal{R}, \rho)$

$$\langle \mathcal{S}(\mathcal{R}, \Sigma) \rangle = \langle L_S(\mathcal{R}, \rho) \rangle$$

ON THE OTHER HAND

$$\langle \mathcal{S}(\mathcal{R}, \Sigma) \rangle =$$

$$= \text{Tr}_{\mathcal{R}} \left(P_{\text{exp}} \int_{\mathcal{S}} \frac{\pi R \varphi}{S} + A + \pi R S \bar{\varphi} \right)$$

$$= \text{Tr}_{\mathcal{R}} \left(\text{Hol}_{\rho}(A) \right)$$

(D.) STRATEGY FOR
COMPUTING FRAMED BPS DEG'S

- $A = \frac{R}{S} \varphi + A + RS \bar{\varphi}$

IS FLAT

- $\text{Tr}_R(\text{Hol}_p A)$ IS HOLOMORPHIC ON \mathcal{M}^J

- \Rightarrow NEED DARBOUX EXPANSION:

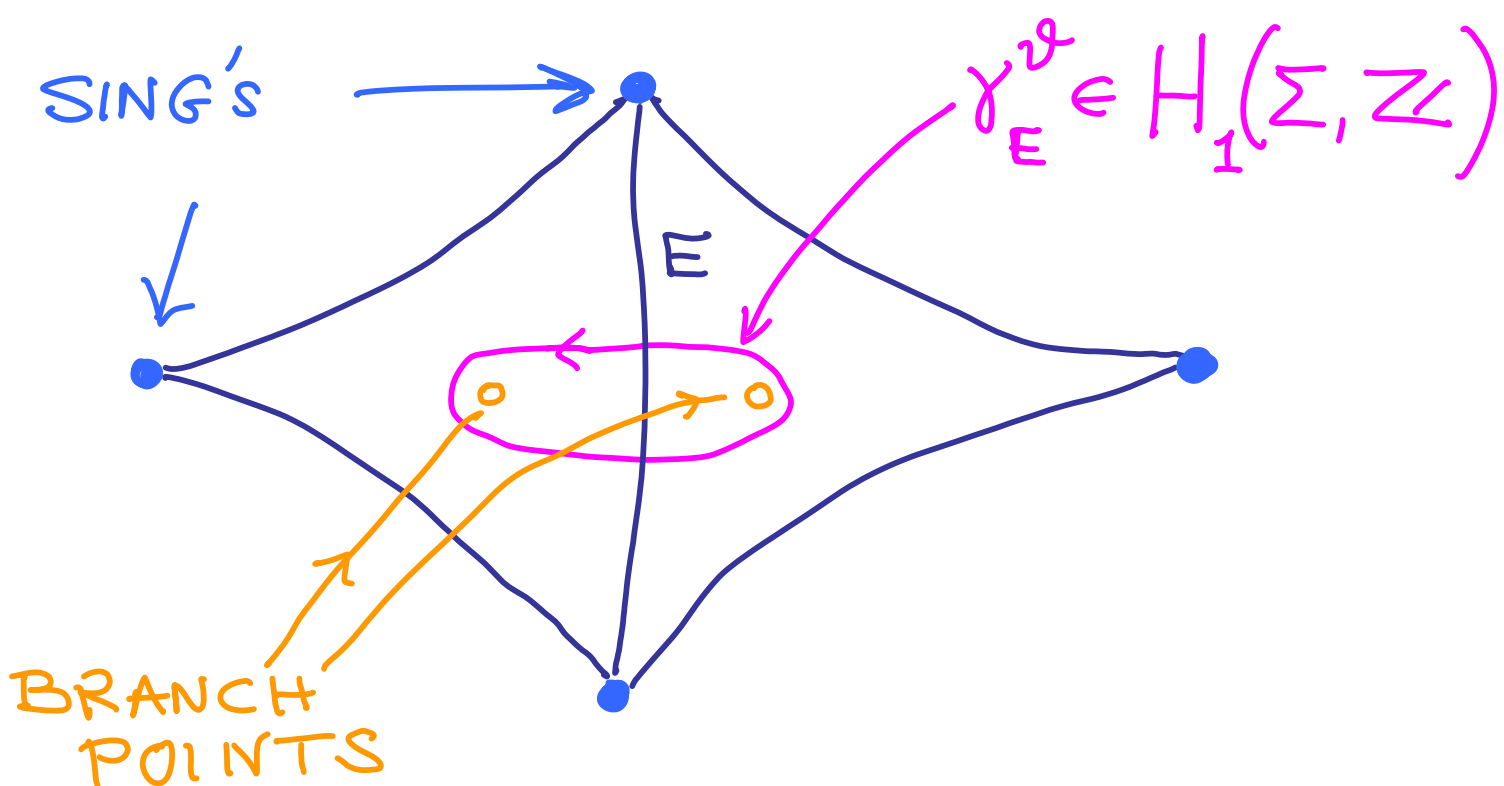
$$\text{Tr}_R(\text{Hol}_p A) = \sum_{\gamma} \underline{\underline{\Omega}}(L(p), \gamma) \gamma_{\gamma}$$

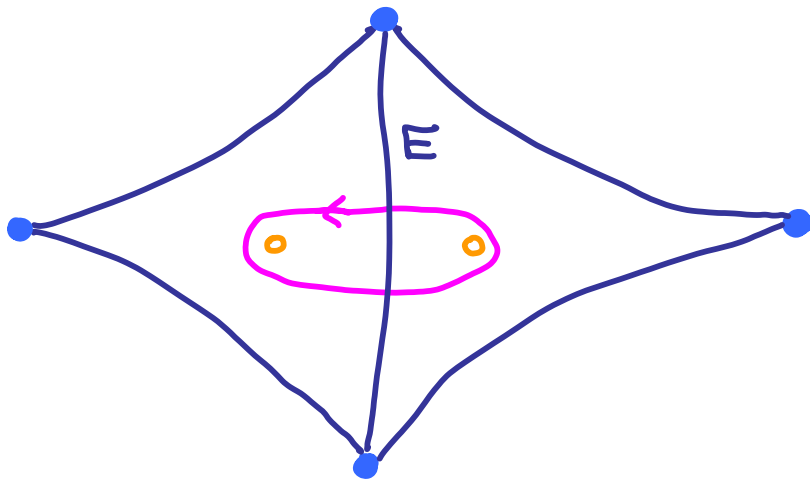
(E.) A_1 -THEORIES: HITCHIN SYSTEMS ON C WITH $\mathfrak{g} = \mathfrak{sl}(2)$

- $\Sigma \rightarrow C$ DOUBLE COVER

- FOR ANGLE ϑ DEFINE A DISTINGUISHED (DECORATED) TRIANGULATION OF C :

WKB CURVE: $\langle \lambda, \partial_t \rangle = e^{i\vartheta}$





$\gamma_{\delta_E} =$ FOCK-GONCHAROV
 COORDINATE X_E W.R.T.
 WKB TRIANGULATION

$\Rightarrow \exists$ ALGORITHMS FOR

COMPUTING $\Omega(\gamma)$ AND $\underline{\underline{\Omega}}(\gamma)$

PURELY IN TERMS OF THE
 COMBINATORICS OF THE TRIANGULATION.

(ONLY EXAMPLE WE HAVE
 WHERE BPS SPECTRUM IS COMPUTABLE
 BY AN ALGORITHM)

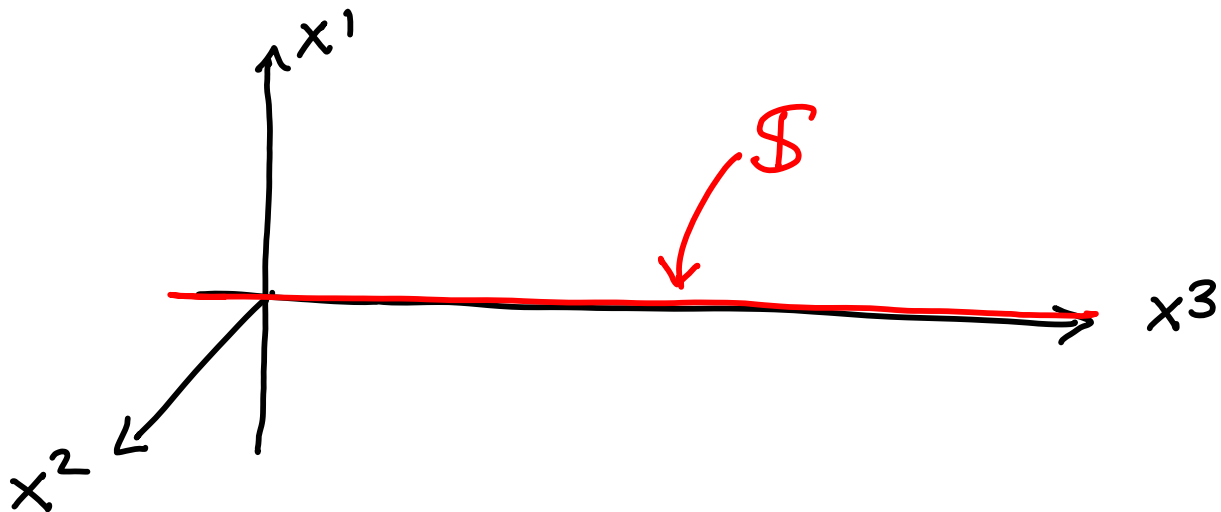
VIII. SURFACE DEFECTS $\frac{1}{2}$ 2D/4D WCF

[MOTIVATION: PAPER BY GAIOTTO, 0911.1316]

(A.) UV DEFINITION

\mathcal{S} SUPPORTED ON $S \subset M^4$

WE TAKE: $x^1 = x^2 = 0$ IN $\mathbb{R}^{1,3}$



AS BEFORE, DEFINE THE DEFECT BY
SCFML BDRY COND. FOR UV
THEORY \mathcal{A}_* ON $AdS_3 \times S^1$

- PRESERVE SUPERALGEBRA

EVEN: $so(2,2) \oplus so(2)_{12} \oplus (U(1) \oplus U(1))_R$

ODD: $\{ Q_\alpha^A, \bar{Q}_{\dot{\alpha}A} \}_{\alpha=A}$

- POINCARÉ' SUBALGEBRA

$$\{ P_0, P_3, M_{12}, F_A, F_V \}$$

\oplus

$$\{ Q_1', Q_2', \bar{Q}_{i1}, \bar{Q}_{i2} \}$$

IS ISOMORPHIC TO $D=2$ $(2,2)$ SUSY

(WITHOUT SETTING $P_1=P_2=0$!!)

$$\left(\begin{matrix} N=2 \\ VM \end{matrix} \right) \supset D=2 (2,2) \text{ TWISTED CHIRAL MULTIPLET}$$

$$\gamma = \varphi + \theta \cdot \lambda + \theta^2 \left(F_{03} - i(F_{12} - D_{12}) \right)$$

(B.) EXAMPLES

①

- CHOOSE $D=2$ $(2,2)$ THEORY T_{2d} ON S WITH G-FLAVOR SYMMETRY
- CHOOSE $D=4$ $N=2$ THEORY T_{4d} WITH G-GAUGE SYMMETRY
- USING RESTRICTION OF TWISTED CHIRAL MULTIPLY, GAUGE THE G-FLAVOR SYMMETRY.

②

- USING $D=4$ THEORY T_{4d} WITH G-GAUGE SYMMETRY, REDUCE THE STRUCTURE GROUP ALONG $S \subset M_4$ TO $T \subset G$ AND INTRODUCE:

$$S = \int_{M_4} d^4x d^4\theta \mathcal{F} + \int_S dx^0 dx^3 d^2\tilde{\theta} W(\gamma) + \text{c.c.}$$

$$\gamma \in \mathfrak{t} = \text{Lie}(T)$$

(C.) IR DESCRIPTION

- ASSUME \mathcal{S} HAS A FINITE SET OF MASSIVE VACUA $i \in \mathcal{V}$
- e.g. $T_{2d} = LG$ MODEL WITH MORSE CRITICAL POINTS
- VACUA FOR THE 2D-4D SYSTEM ARE DESCRIBED BY (u, i) AND TOGETHER FORM A FINITE COVER:

$$\mathcal{V} \longrightarrow \mathcal{B}_{\mathcal{S}} \\ \downarrow \\ \mathcal{B}$$

IR LAGRANGIAN

- 4D: UV \rightarrow IR

$$\mathcal{F} \rightarrow \mathcal{F}^{\text{eff}} \left(V^I = a^I + \dots \right)$$

\uparrow
4D ABELIAN VM'S
IN SOME DUALITY FRAME

-
- 2D+4D: UV \rightarrow IR

$$S^{\text{eff}} = \int d^4x d^4\theta \mathcal{F}^{\text{eff}}(V^I) + \int dx^0 dx^3 d^2\tilde{y} W^{\text{eff}}(\gamma^I)$$

\mathcal{F}^{eff}

DEPENDS ON $u \in \mathcal{B}$

W^{eff}

DEPENDS ON $(u, i) \in \mathcal{B}_{\mathcal{S}}$

$\mathcal{F}^{\text{eff}}, W^{\text{eff}}$

ONLY LOCALLY DEFINED

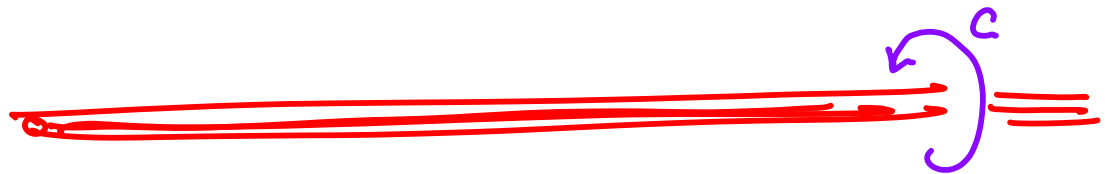
(D.) GUKOV - WITTEN PARAMETERS

$$L_I = \frac{\partial W^{\text{eff}}}{\partial a^I} = \eta_I + \tau_{IJ} \alpha^J$$

PHYSICAL INTERPRETATION:

SOLENOID

FLUX
= δ_i
FOR
VAC = ?



$$F = dA = 0$$

$$\text{BUT } \oint_c A \in \bar{V} = \Gamma \otimes \mathbb{R}$$

CHOOSE DUALITY FRAME

$$\oint_c A = \begin{pmatrix} \eta_I \\ \alpha^I \end{pmatrix}$$

N.B. AHARONOV-BOHM PHASE OF A PROBE PARTICLE

$$\gamma_t = \begin{pmatrix} p_t \\ q_t \end{pmatrix}$$

$$\text{IS } \exp 2\pi i \left(p_t^I \eta_I - q_t^I \alpha^I \right)$$

- GAUGE TRANS
- MONODROMY \Rightarrow SHIFTS

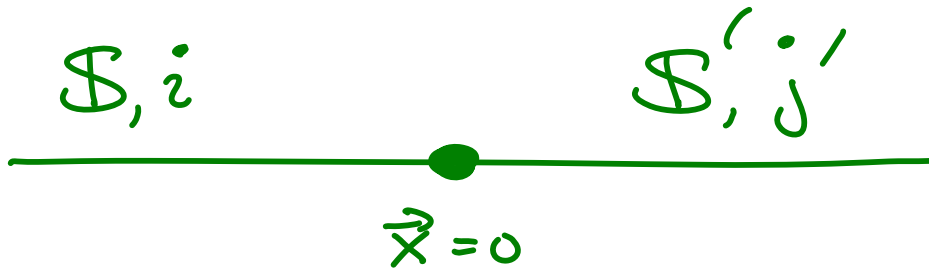
$$\begin{pmatrix} \eta \\ \alpha \end{pmatrix} \rightarrow \begin{pmatrix} \eta \\ \alpha \end{pmatrix} + \gamma \quad \gamma \in \Gamma$$

\Rightarrow G-W PARAMETERS LIVE
IN A Γ -TORSOR: Γ_i

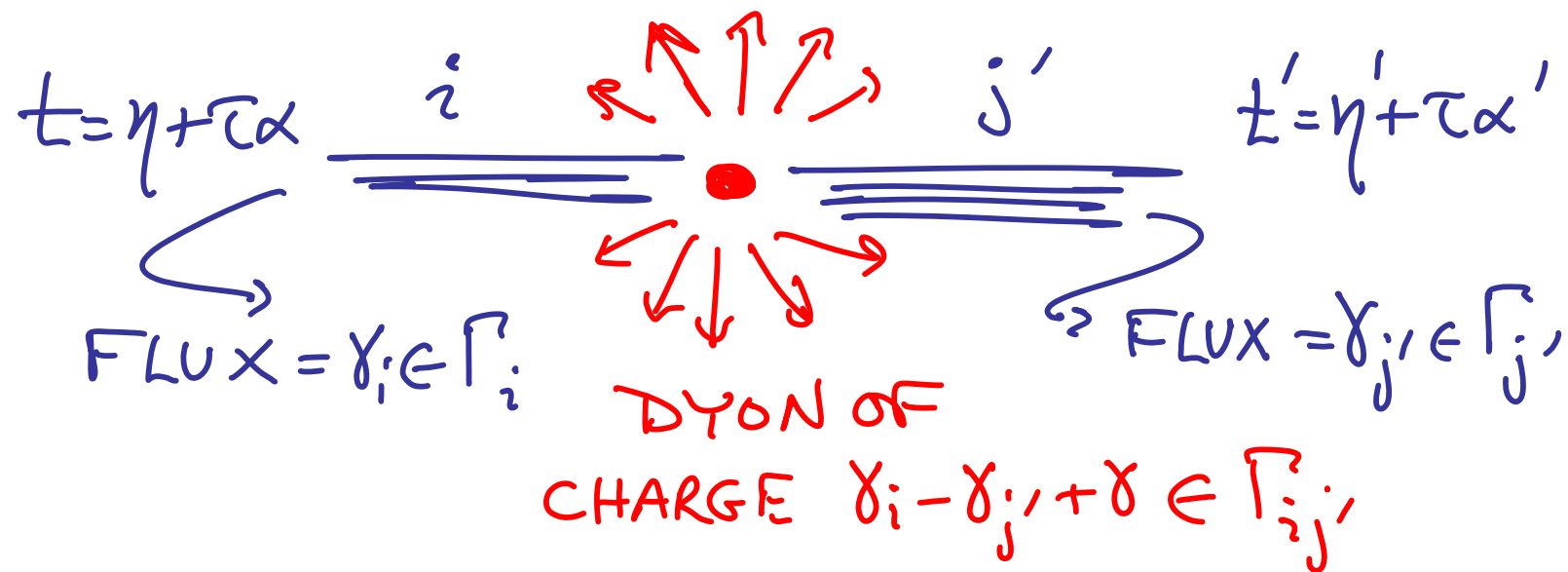
$$\gamma_i = \text{FLUX IN SOLENOID}$$

(E.) INCLUDING LINE DEFECTS

LINE DEFECTS = DOMAIN WALLS



IR: DYON EMBEDDED IN
A SOLENOID:



DYON CHARGE LIVES IN Γ -TORSOR

$\Gamma_{ij'}$

(F.) BPS DEGENERACIES

- $\Omega(\gamma; u)$ AS BEFORE
- $\mu(\gamma_{ij})$: # 2D SOLITONS BETWEEN VAC $i \& j$ OF 4D CHARGE $\gamma_{ij} \in \Gamma_{ij}$
- $\omega(\gamma, \gamma_i)$: # 2D PARTICLES IN VACUUM i OF 4D CHARGE γ

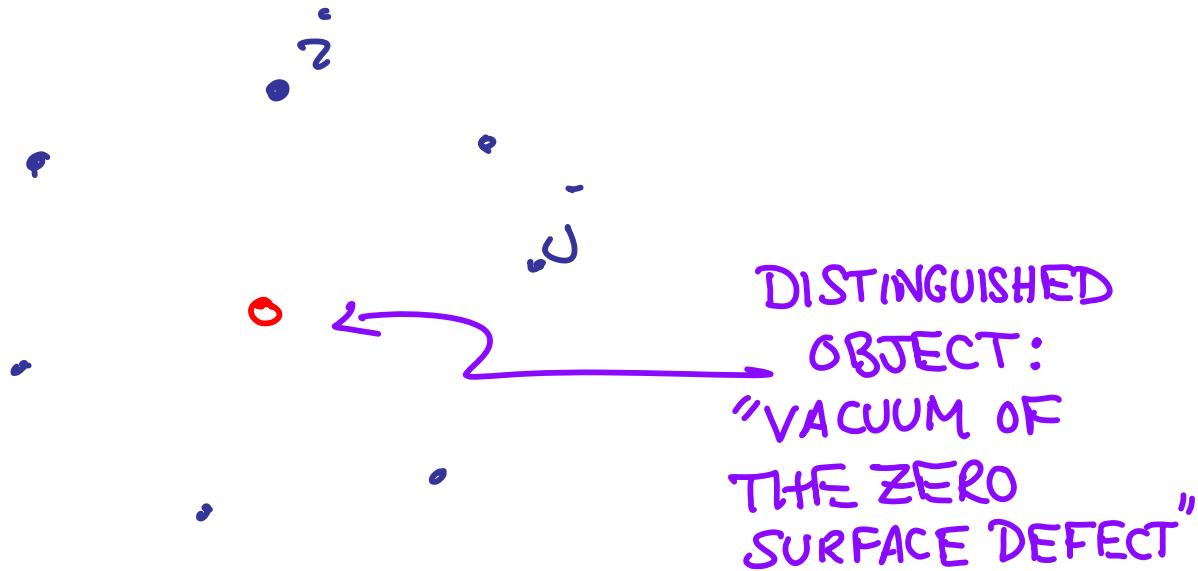
$\sum \omega(\gamma, \gamma_i)$ IS SUBTLE TO DEFINE BECAUSE THESE STATES ARE AT THRESHOLD

DEPENDS AFFINE-LINEARLY ON FLUX γ_i

(G.) FORMAL STATEMENT OF 2D4D WCF

4 PIECES OF DATA

① GROUPOID OF VACUA \mathcal{V} :

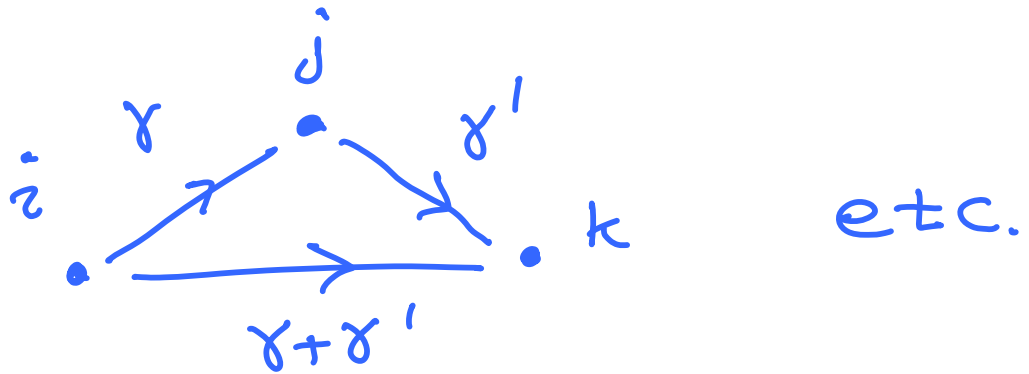


MORPHISMS:
$$\begin{array}{ccc} \dot{z} & \gamma & \dot{j} \\ \bullet & \longrightarrow & \bullet \end{array}$$

DRAW AN ARROW FOR EACH $\gamma \in \Gamma$ TO GET $\text{Hom}(i, j) = \Gamma$ -torsor.

$$\text{Hom}(0, 0) = \Gamma, \quad \text{Hom}(i, 0) = \Gamma_i, \quad \dots$$

COMPOSITION OF MORPHISMS



LET $a \in \Gamma, \Gamma_i, \Gamma_{ij}$

DENOTE COMPOSITION $a+b$
WHEN DEFINED

② CENTRAL CHARGE $Z \in \text{Hom}(V, \mathbb{C})$

$$Z(a) + Z(b) = Z(a+b)$$

WHEN $a+b$ IS DEFINED

③ BPS DATA:

$$\Omega(\gamma) \in \mathbb{Z}$$

$$\mu(\gamma_{ij}) \in \mathbb{Z}$$

$$\omega(\gamma, \gamma_a) \in \mathbb{Z}$$

$$\omega(\gamma, \gamma_a + \gamma') = \omega(\gamma, \gamma_a) + \Omega(\gamma) \langle \gamma, \gamma' \rangle$$

PIECEWISE CONSTANT IN
STABILITY DATA \mathbb{Z}

④ TWISTING FUNCTION

$$\sigma(a, b) \in \mathbb{Z}_2 \text{ WHEN } a+b \text{ DEFINED}$$

$$\sigma(a, b) \sigma(a+b, c) = \sigma(a, b+c) \sigma(b, c)$$

3 DEFINITIONS

① DEF: A BPS RAY IS A
RAY IN CPLX PLANE:

- $\mathbb{R} \cdot Z(\gamma)$ IF $\omega(\gamma, \cdot) \neq 0$
- $\mathbb{R} \cdot Z(\gamma_{ij})$ IF $\mu(\gamma_{ij}) \neq 0$

② DEFINE THE (TWISTED)
GROUPOID ALGEBRA $\mathbb{C}[W]$:

$$X_a X_b = \begin{cases} \sigma(a,b) X_{a+b} & \text{IF } a+b \\ & \text{COMPOSABLE} \\ 0 & \text{ELSE} \end{cases}$$

3. DEFINE TWO AUTOMORPHISMS OF $\mathbb{C}[V]$:

• CV-LIKE: $S_{\gamma_{ij}}^{\mu}$

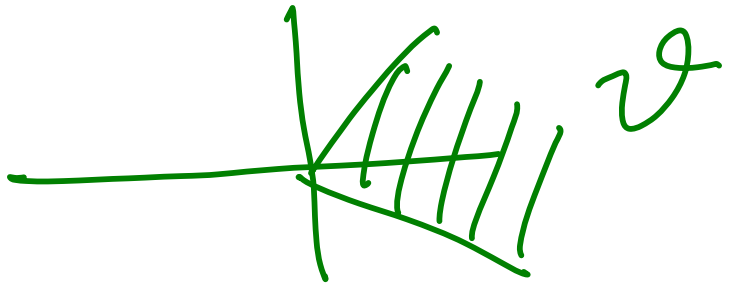
$$X_a \rightarrow (1 - \mu(\gamma_{ij})X_{\gamma_{ij}}) \cdot X_a \cdot (1 + \mu(\gamma_{ij})X_{\gamma_{ij}})$$

• KS-LIKE K_{γ}^{ω} :

$$X_a \rightarrow (1 - X_{\gamma})^{\omega(\gamma, a)} X_a$$

2D/4D WALL-CROSSING FORMULA

FOR CONVEX
SECTOR



$$A(\mathcal{V}) = : \prod_{Z(\gamma_{ij}) \in \mathcal{V}} S_{\gamma_{ij}}^M \prod_{Z(\gamma) \in \mathcal{V}} \prod K_{\gamma}^{\omega} :$$

ORDERED BY PHASE OF Z

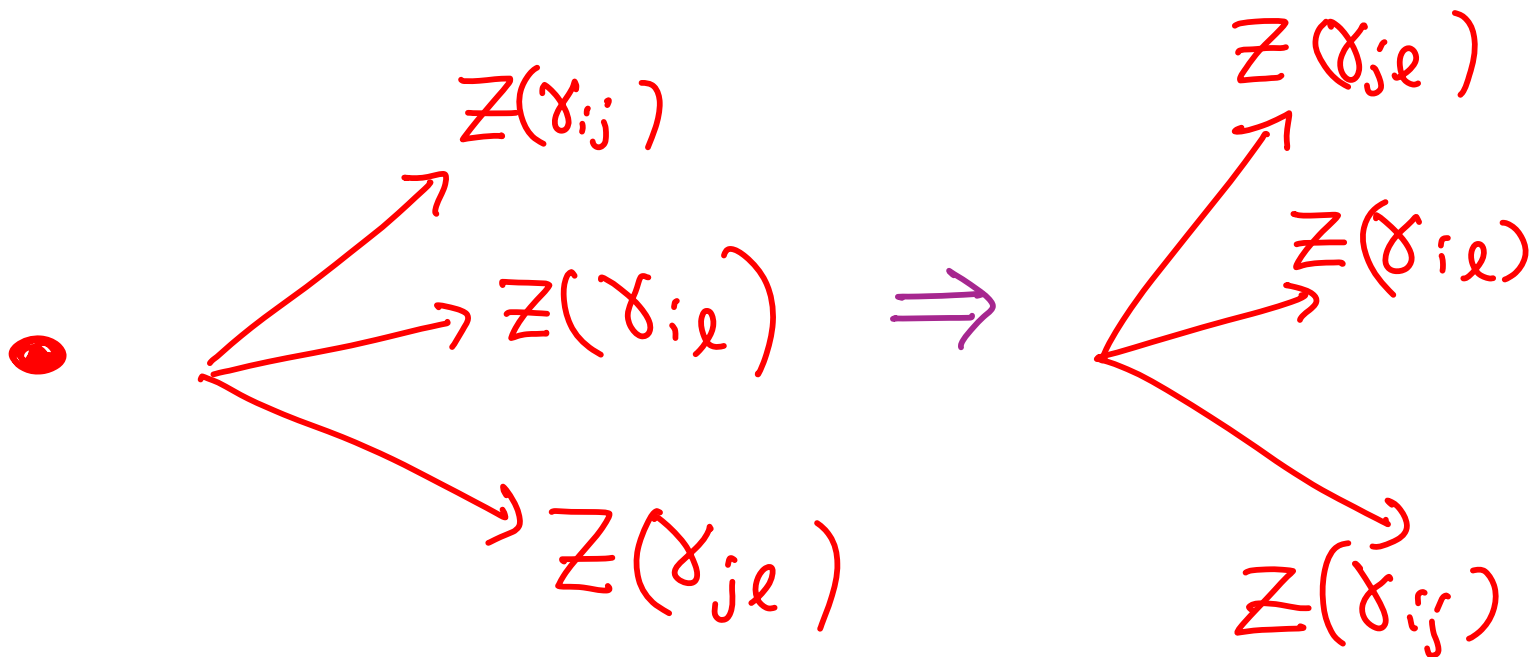
W.C.F.: $A(\mathcal{V})$ IS

CONSTANT SO LONG AS

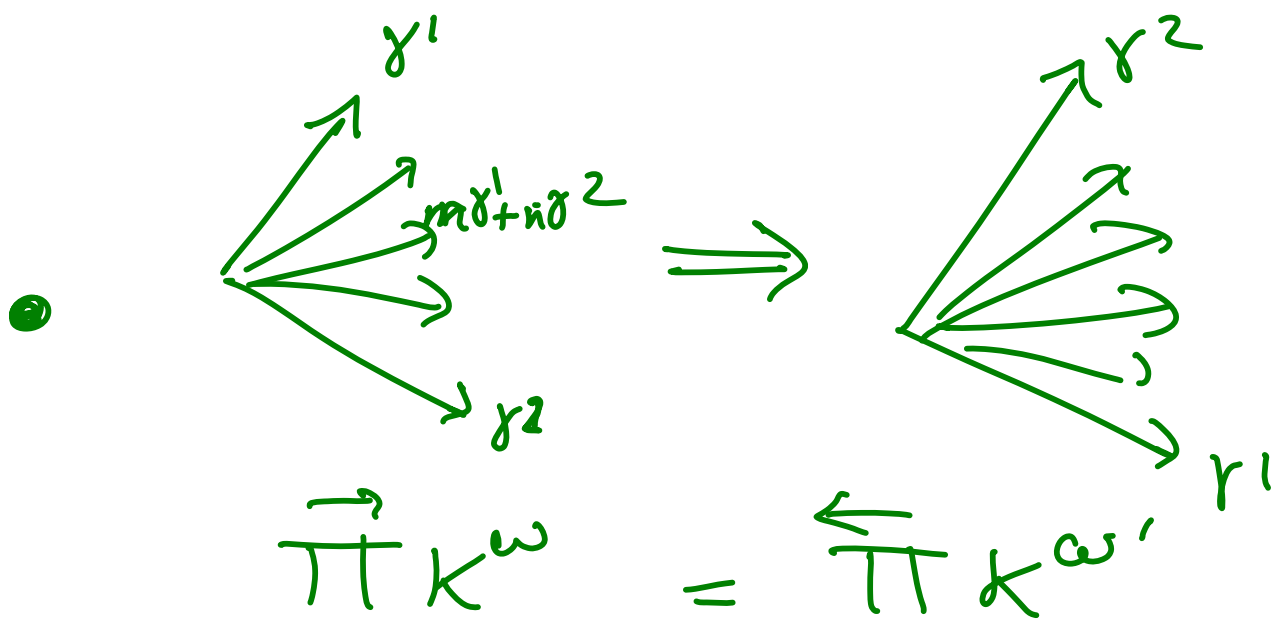
NO BPS LINE ENTERS/LEAVES
THE SECTOR \mathcal{V} .

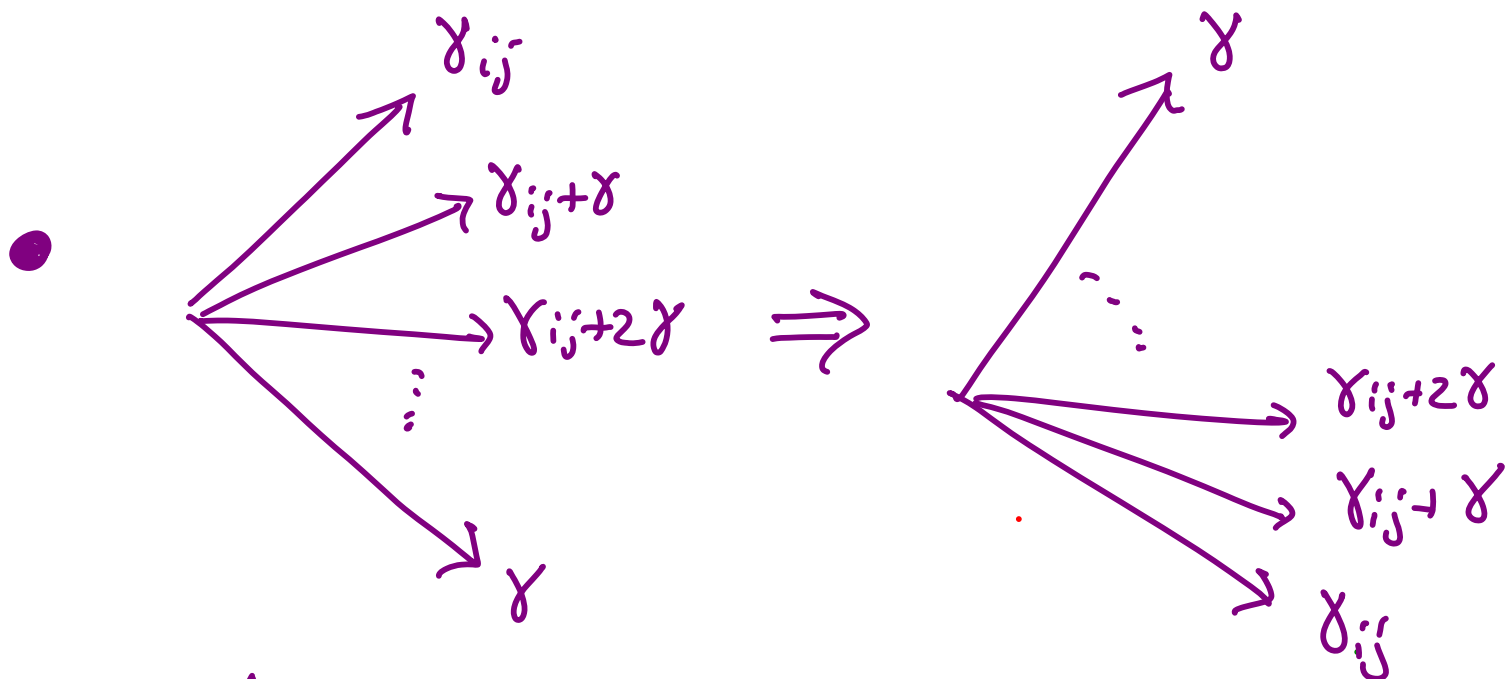
CV + KS WCF ARE
SPECIAL CASES

TYPES OF WALL-CROSSING

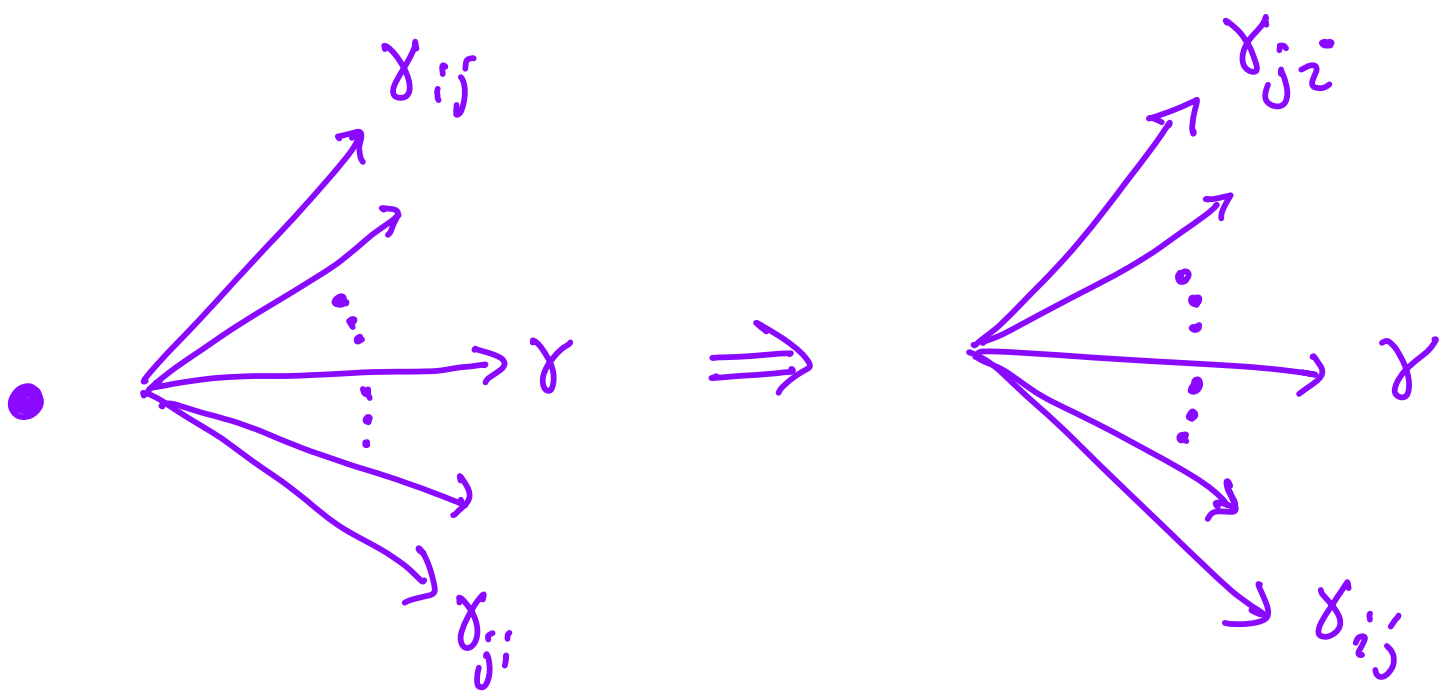


$$S_{\gamma_{ij}}^{\mu} S_{\gamma_{ie}}^{\mu} S_{\gamma_{je}}^{\mu} = S_{\gamma_{je}}^{\mu'} S_{\gamma_{ie}}^{\mu'} S_{\gamma_{ij}}^{\mu'}$$





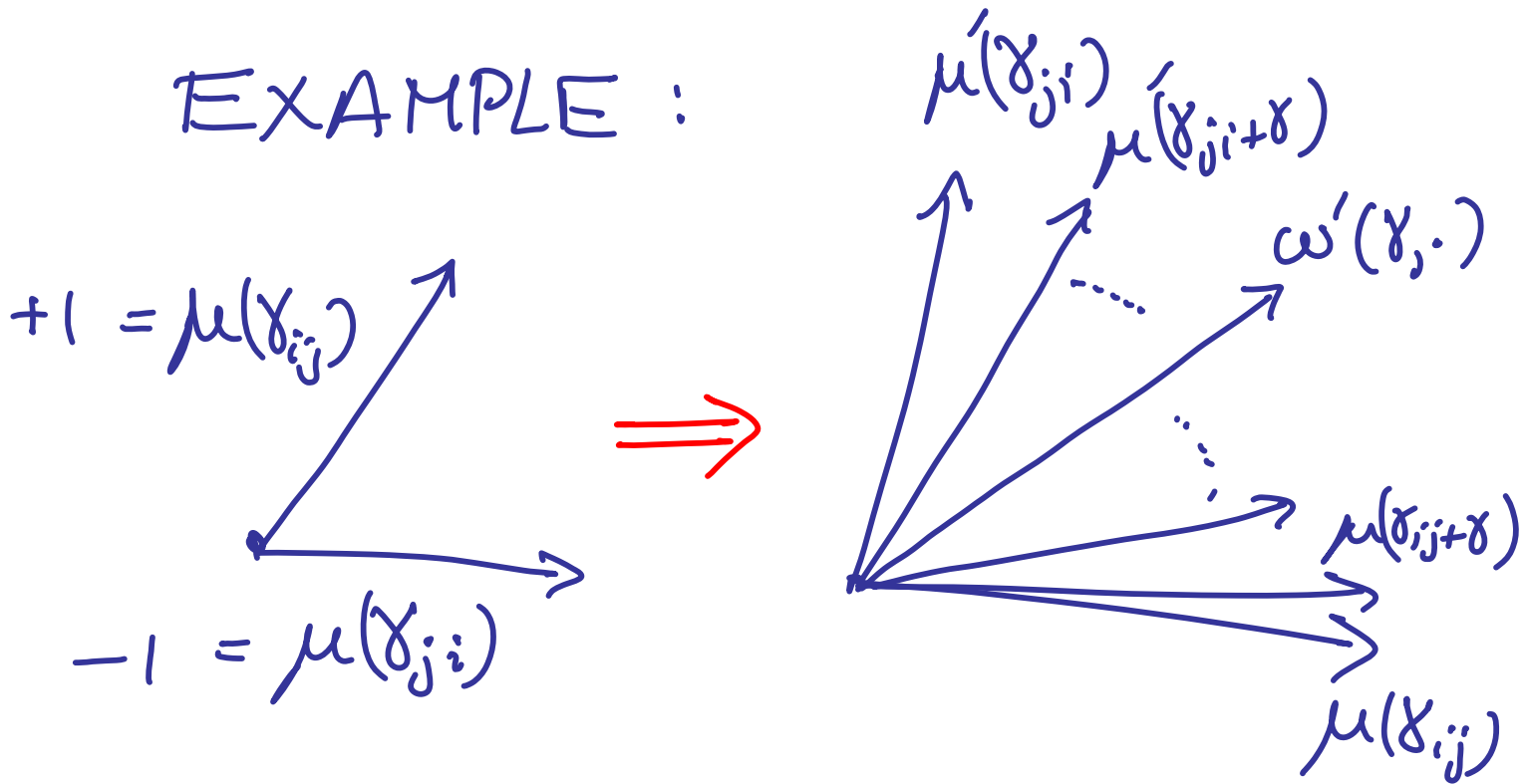
$$K_{\gamma}^{\omega} \prod S_{\gamma_{ij+n\delta}}^{\mu} = \prod S_{\gamma_{ij+n\delta}}^{\mu'} K_{\gamma}^{\omega'}$$



$$\prod S_{\gamma_{ij+n\delta}}^{\mu} K_{\gamma}^{\omega} \prod S_{\gamma_{ij+n\delta}}^{\mu} = \prod S_{\gamma_{ji+n\delta}}^{\mu'} K_{\gamma}^{\omega'} \prod S_{\gamma_{ji+n\delta}}^{\mu'}$$

THM: MIXED WCF CAN BE SOLVED EXPLICITLY

EXAMPLE :

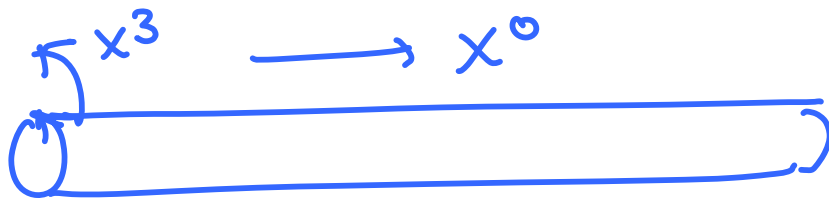


$$S_{\gamma_{ij}}^{\mu} S_{\gamma_{ji}}^{\mu} = \prod_{\nearrow} S_{\gamma_{ji+n\delta}}^{\mu'} K_{\gamma}^{\omega'} \prod_{\searrow} S_{\gamma_{ij+n\delta}}^{\mu'}$$

ENCODES \mathbb{CP}^1 COUPLED TO TWISTED MASS PARAMETER.

(CECOTTI ; DOREY)

(H.) 3D GEOMETRY



SURFACE DEFECT IN 4D THEORY



LINE DEFECT IN 3D THEORY

$$\exp i \int t_{\mathbf{I}} (F_{03}^{\mathbf{I}} - i F_{12}^{\mathbf{I}}) + \text{c.c.} \, dx^0 dx^3$$



$$\exp i \int dx^0 \left(\eta_{\mathbf{I}} d\theta_{ee}^{\mathbf{I}} - \alpha^{\mathbf{I}} d\theta_{mg\mathbf{I}} \right)$$

$$d\theta_{ee}^{\mathbf{I}} = \oint_{S'_R} F^{\mathbf{I}} \quad d\theta_{mg,\mathbf{I}} = \oint_{S'_R} G_{\mathbf{I}}$$

$\eta_I d\theta_{el}^I - \alpha^I d\theta_{mg,I}$ IS A

LOCALLY DEFINED 1-FORM ON \mathcal{M}

CROSSING PATCHES WE FIND

A CONNECTION A^{sf} ON

$$(\mathcal{L}_S)_i \rightarrow \mathcal{M}$$

WHEN \mathcal{M} HAS THE
SEMI FLAT HK METRIC

A_i^{sf} IS HYPER-HOLOMORPHIC

(F_i^{sf} IS TYPE (1,1) IN ALL
COMPLEX STRUCTURES)

QUANTUM CORRECTIONS

$$\begin{array}{ccc} \mathcal{L}_S & \rightarrow & \mathcal{M}_S \xrightarrow{\pi} \mathcal{M} \\ & & \downarrow \qquad \downarrow \\ & & \mathcal{B}_S \rightarrow \mathcal{B} \end{array}$$

$$V_S = \pi_* (\mathcal{L}_S) = \bigoplus_{i \in V} (\mathcal{L}_S)_i$$

HAS A CONNECTION \oplus

3D LINE DEFECT $\text{Tr } P \exp \int dx^0 \varphi^* \oplus$

SUSY $\Rightarrow \oplus$ IS HYPERHOLOMORPHIC (THOMPSON)

CONSTRUCT \oplus VIA ITS HOLOMORPHIC SECTIONS

$$\psi_{\gamma_i} \in H^0(V_S)$$

$$\begin{array}{c}
 \mathcal{U}_i \xrightarrow{\quad} (\mathcal{L}_S)_i \\
 \xrightarrow{\quad} \\
 \mathcal{U} \xrightarrow{\quad} \\
 \downarrow
 \end{array}$$

Split: $\gamma_{\mathcal{X}_i} = g_i \gamma_{\mathcal{X}_i}$

$$\gamma_{\mathcal{X}_i} \in H^0((\mathcal{L}_S)_i, \mathcal{U})$$

$$g_i \in \text{Hom}((\mathcal{L}_S)_i, \bigoplus_j (\mathcal{L}_S)_j)$$

$$\gamma_{\mathcal{X}_i} = \gamma_{\mathcal{X}_i}^{\text{st}} \exp \left[- \sum_{\mathcal{X}} \omega(\mathcal{X}, \mathcal{X}_i) K_{\mathcal{X}} * \log(1 - \mathcal{X}_{\mathcal{X}}) \right]$$

$$g_i = g_i^0 - \sum_{\substack{j \neq i \\ \mathcal{X}_{ji}}} \mu(\mathcal{X}_{ji}) K_{\mathcal{X}_{ji}} * (g_j \gamma_{\mathcal{X}_{ji}})$$

(I.) HITCHIN SYSTEMS

$z \in \mathbb{C} \implies \mathcal{S}_z$ SURFACE DEFECT

$V(\mathcal{S}_z) = \text{UNIVERSAL BUNDLE} \rightarrow \mathcal{M}$

VACUA = PREIMAGES $\Sigma \rightarrow \mathbb{C}$ (GAIOTTO)
 $x_i \rightarrow z$

$\Gamma_{ij} = \{ \text{Paths } \gamma_{ij}: x_i \text{ TO } x_j \} / \partial(2\text{-CHAN})$

$$\omega(\gamma, \gamma_{ij}) = \Omega(\gamma; u) \langle \gamma, \gamma_{ij} \rangle$$

γ : WKB Rep.

$\mu(\gamma_{ij}) = \text{SIGNED SUM OF WKB CURVES } x_i \text{ TO } x_j \text{ THROUGH RAM. POINT OF TYPE } (ij)$

IX. SOME OPEN PROBLEMS

1. GENERALIZE A_1 -THEORIES TO HIGHER RANK: QUALITATIVELY NEW PHENOMENA
2. QUANTIZATION OF \mathcal{M} : RELATION TO NEKRASOV-SHATASHVILI \S NEKRASOV-WITTEN?
3. APPLICATIONS TO SUGRA?
4. $W=4$ SCATTERING?
5. NEW MODULAR FUNCTORS.
6. $\overline{\Sigma}$, PSC, \S TOPOLOGY OF MONOPOLE MODULI SPACES.
7. STRONG POSITIVITY \S WORK OF LUSZTIG, FOMIN-ZELEVINSKY, ...
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